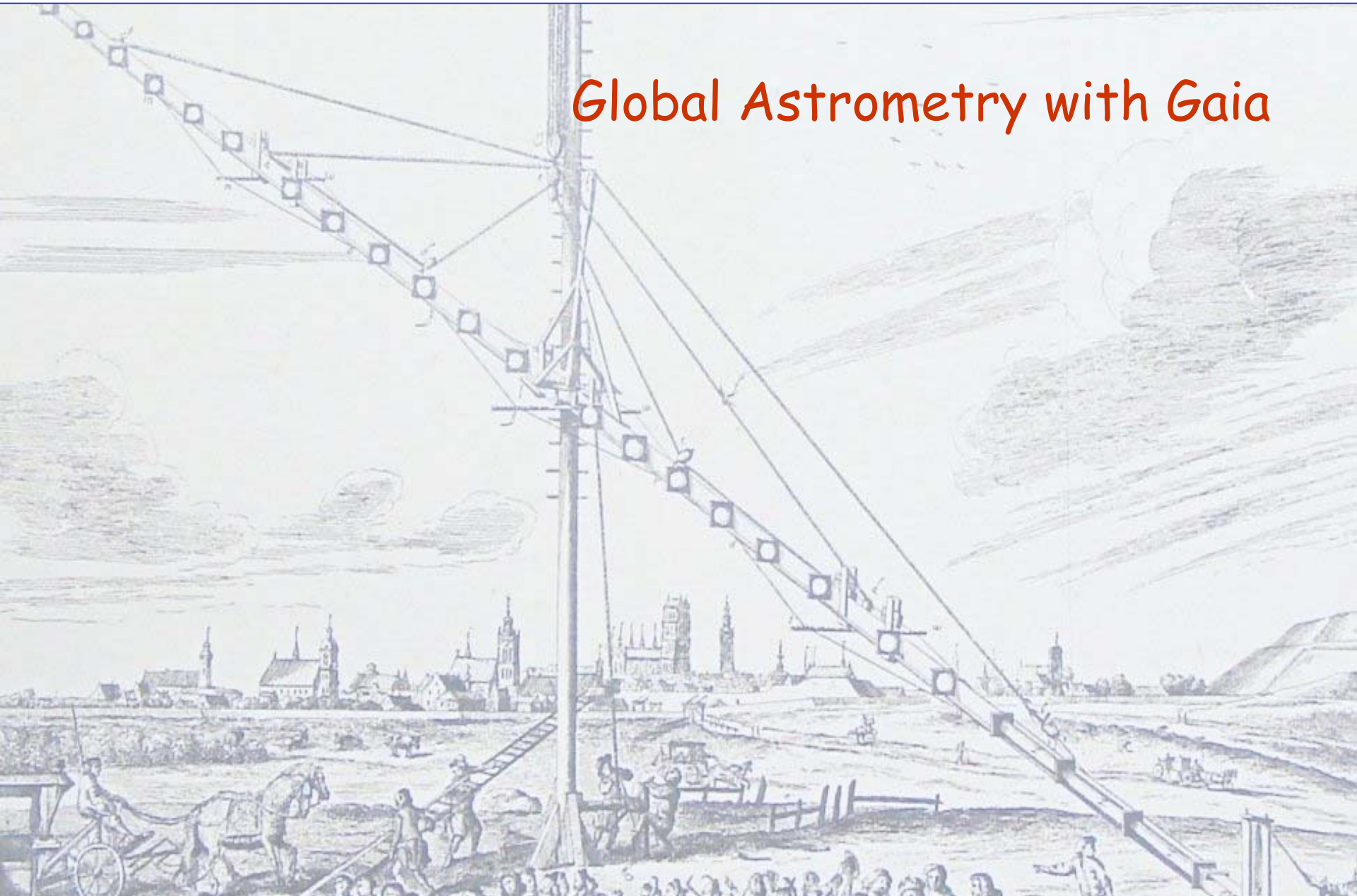


Global Astrometry with Gaia

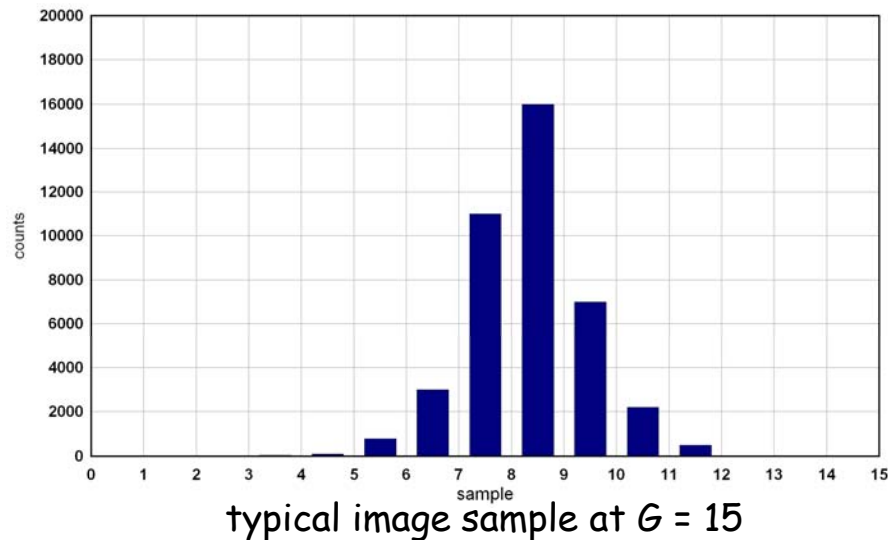


- Small field astrometry can be very accurate, even more than Gaia
 - but it cannot be used to relate fields at large angular distance
- How to eliminate the systematic distortions at large angle ?
 - One needs to have direct and numerous connections between sources at large angular separations
 - in a small field the distortions are not seen, as they are constant over the frame size
 - in a very very large field, or in two **connected** fields, the possible distortions are different in each field and are included in the model
- To produce a full-sky frame, observations of many sources, in many directions must be processed together
 - the astrometry is then global

Assume you are on board of a spinning spaceship to do astrometric observations

- If one knew perfectly the rotational motion of the platform, it is possible to map the sky from local measurements
- If one knew where the stars are, it is possible to monitor the attitude of the spacecraft from local measurements
- With Gaia, one knows **neither** in advance and one determines **both** from local measurements !
 - in fact almost both, 6 frame orientation parameters are free
- Keyword : **Connectivity** between sources + smooth attitude

- Gaia observes transit times of point sources across fiducial lines on the detector
 - there are 9 observations per transit on the astrometric FOV
 - this crossing is that of the image centre relative to CCD
 - the local astrometric centring accuracy is about $200 \mu\text{as}$ at $G = 15$
 - this is $1/300$ of a pixel size along-scan
 - achievable with the $\sim 40,000$ counts over 4.5 s of integration
 - there are on the average 700 such measurements per star



- **Central Problem:**
 - For each of 10^9 observed celestial objects we want to determine six **astrophysical parameters:**
 - Position on celestial sphere: α, δ
 - Parallax (distance): π
 - Proper motion: μ_α, μ_δ
 - Radial velocity: v_R
 - at the μas level (π : $<25\mu\text{as}@V=15$, $<7\mu\text{as}@V<10$)
 - using (in theory) no a priori knowledge of these quantities but derive them from observation data alone in a self-consistent manner

- Each observed object entering the FOV transits 9 AF CCDs
 - 9 elemental observations per object per field transit
 - one observation: a CCD centroid of the image in pixel units
- $\sim 10^9$ objects in total
- In 5 years we will have ~ 80 FOV transits per sources
 - 10^{11} transits
 - 10^{12} elementary data

→ 1 trillion observations

- Need to determine $> 5 \times 1$ billion unknowns
from the 10^{12} observations using an observation model that incorporates
 - **Satellite attitude** $\rightarrow 40 \times 10^6$ unknowns
 - **Calibration parameters** $\rightarrow 1$ million unknowns
 - **Global astrophysical parameters** $\rightarrow \sim 10$
- Could set up a system of equations that solves directly for the unknowns - system is manifold over-determined
- Observation model not too complex
- Problem: Connectivity of the parameters

- The transit time of a given star can be predicted from:

- the star astrometric parameters $\alpha, \delta, \varpi, \mu_\alpha, \mu_\delta$

S

- the attitude of the satellite $q(t) : q_0, q_1, q_2, q_3$

A

- the orbit of Gaia in barycentric frame $\mathbf{r}(t), \mathbf{v}(t)$

- the imaging properties of the telescopes, the basic angle between the two fields, a scale factor to link length and angles, the geometry of the CCD mosaic

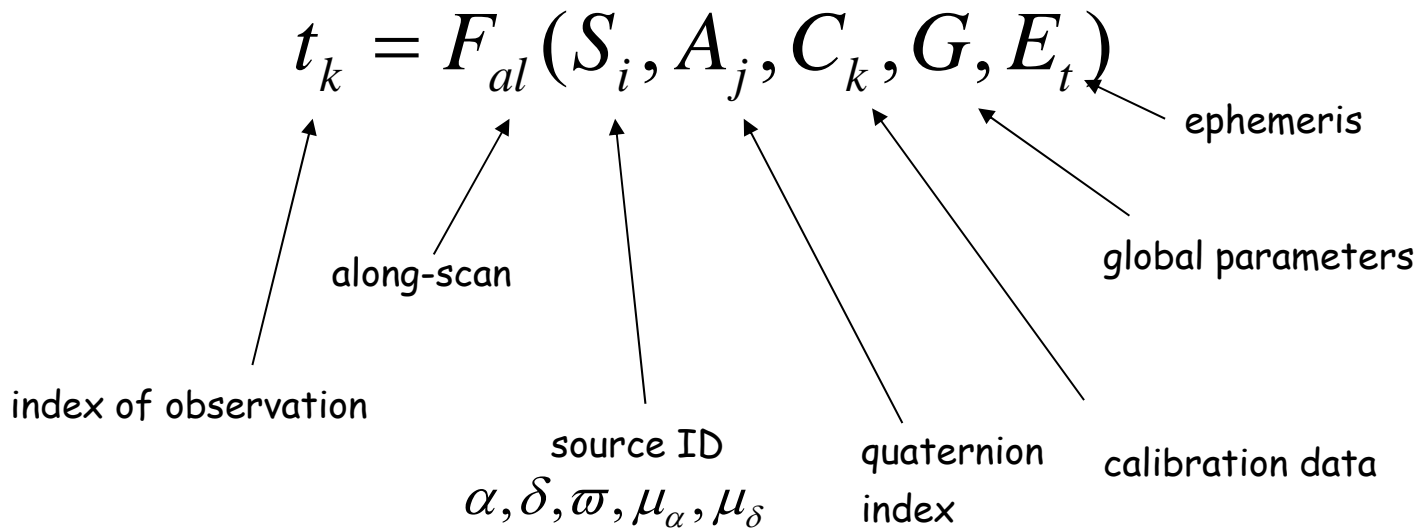
C

- the relationship between on-board clock ticking and TCB

- few general parameters applicable to all stars

G

- Basic direct model



- Example (for a minor planet)

1887.6355353937	2855 Bastian	1	18.77	0.01609	3	571				
490.791788164049	42404410497373876		490.791928426528			42404422616052008		2015-03-03T15:15:10.258027911		
5.093826158	-0.495660889	291.8547404	-28.3992770	0.38392	0.03785	-5.20	13.17	109.66	254.9	

- Linearization about the provisional values

$$R = t_{obs} - t_{comp} = \frac{\partial F_{al}}{\partial S_i} \Delta S_i + \frac{\partial F_{al}}{\partial A_j} \Delta A_j + \frac{\partial F_{al}}{\partial C_k} \Delta C_k + \frac{\partial F_{al}}{\partial G} \Delta G + \varepsilon$$

- One solves for a subset of well-behaved stars
 - several 10^7 , up to 100 millions
- Typically at mission end
 - 500 millions unknowns for stars
 - 20 to 40 millions for attitude
 - 1 million for instrument
 - 100 for general parameters
- Global problem with dense interconnections
 - direct solution not feasible with current means

- Key: relax the connectivity issue

	known		unknown	
Full problem:	R	=	$S + A + C + G + \varepsilon$	

If one knows attitude, instrument ... then	$R - A - C - G$	=	$S + \varepsilon$	10 ⁸ small problems
--	-----------------	---	-------------------	--------------------------------

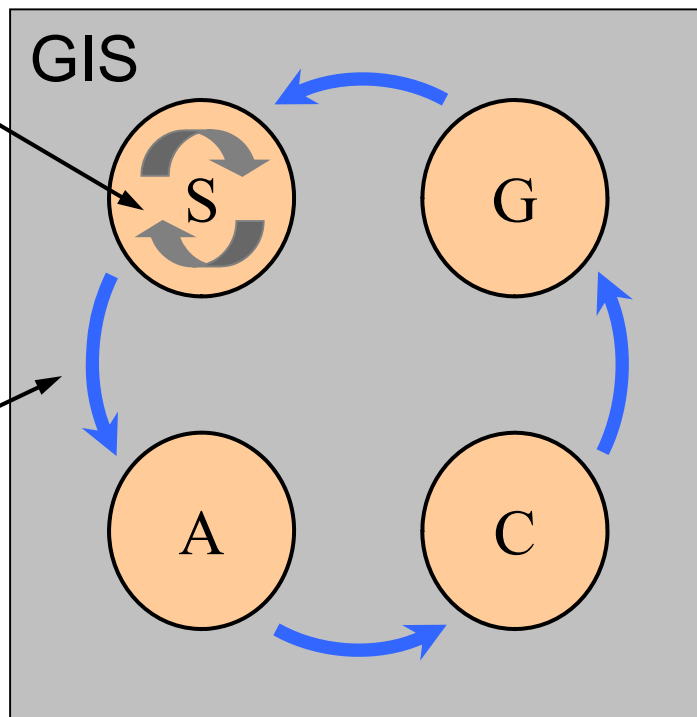
If one knows stars, instrument ... then	$R - S - C - G$	=	$A + \varepsilon$	10 ⁷ small problems
---	-----------------	---	-------------------	--------------------------------

	$R - S - A - G$	=	$C + \varepsilon$	
--	-----------------	---	-------------------	--

	$R - S - A - C$	=	$G + \varepsilon$	
--	-----------------	---	-------------------	--

inner iteration
(non-linearity,
outliers)

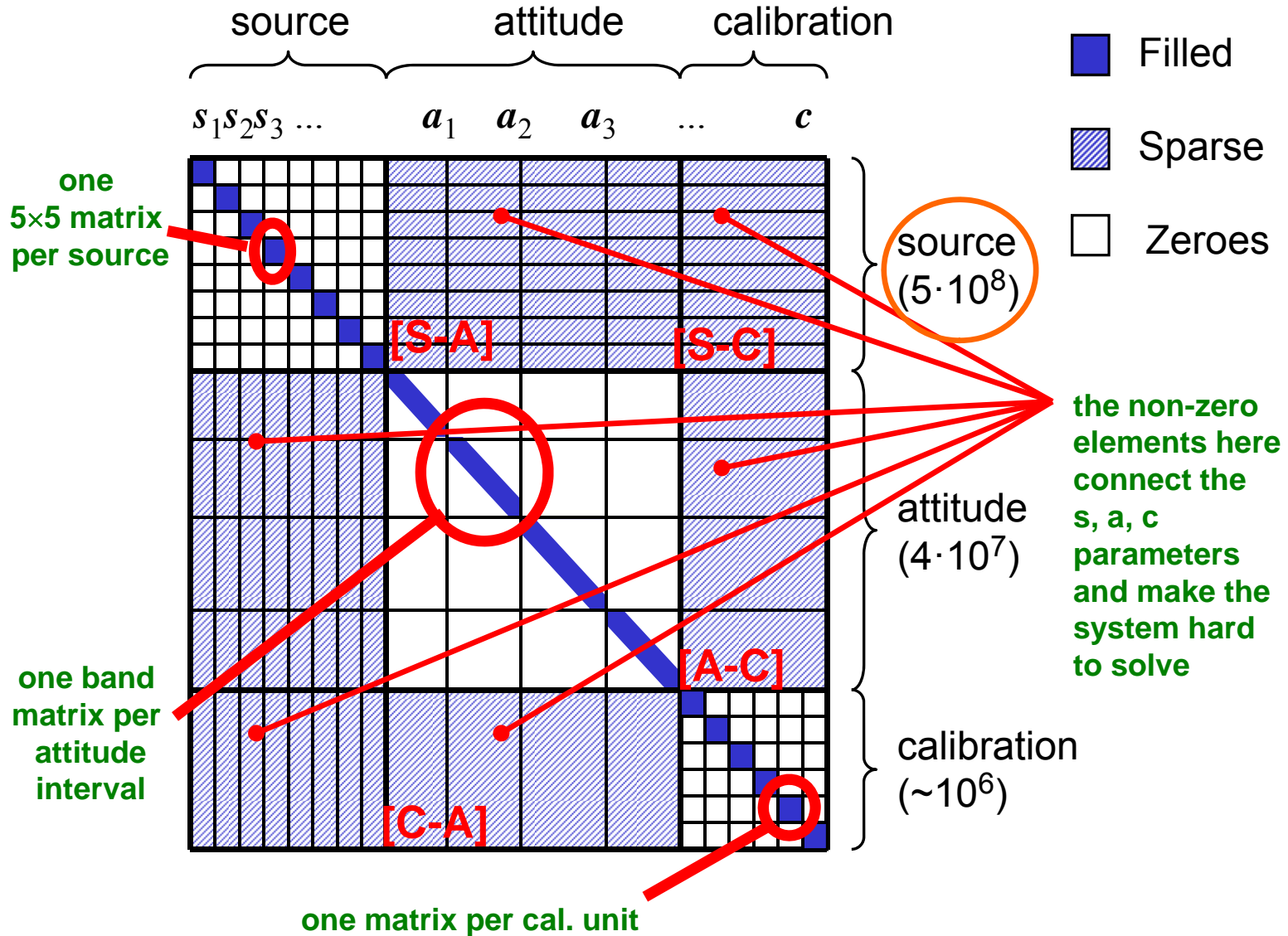
GIS iteration
(S-A-C-G
cross-terms)



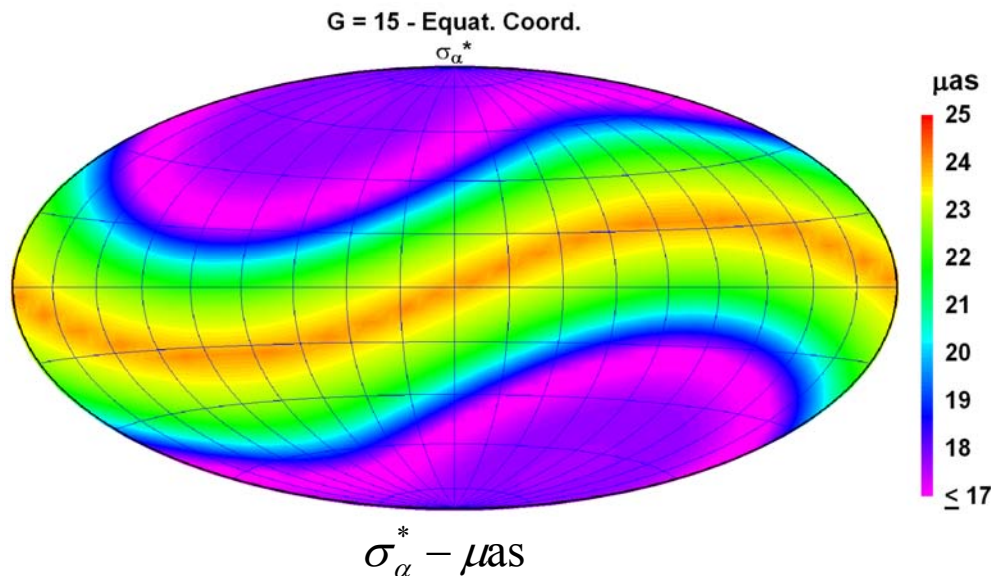
S = Sources
A = Attitude
C = Calibration
G = System, relativity

outer iteration
(interaction
with all other
processes)
Every 6 months

- new data (IDT)
- Improved old data (IDU)
- improved selection of primaries

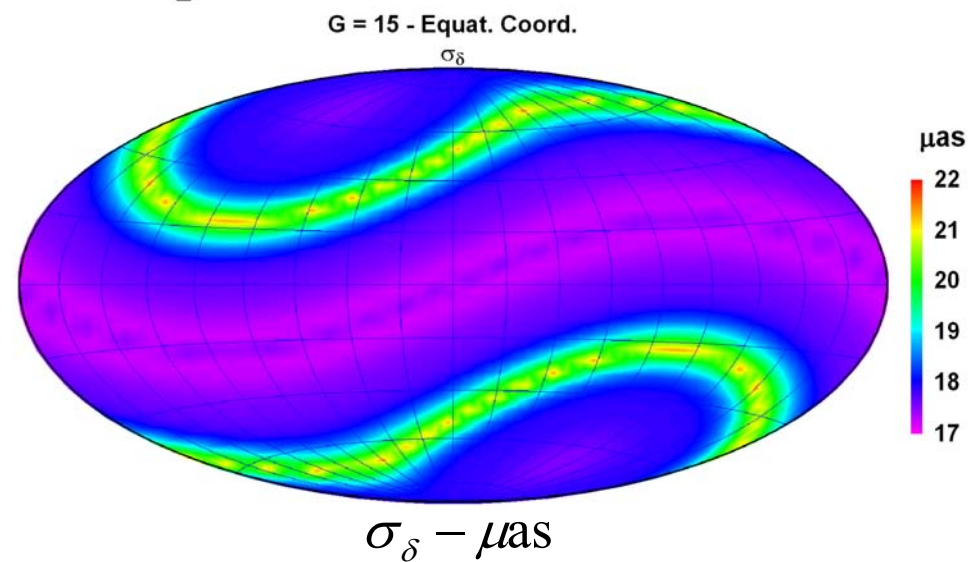


- Plots for $G = 15$, but scalable to other magnitudes



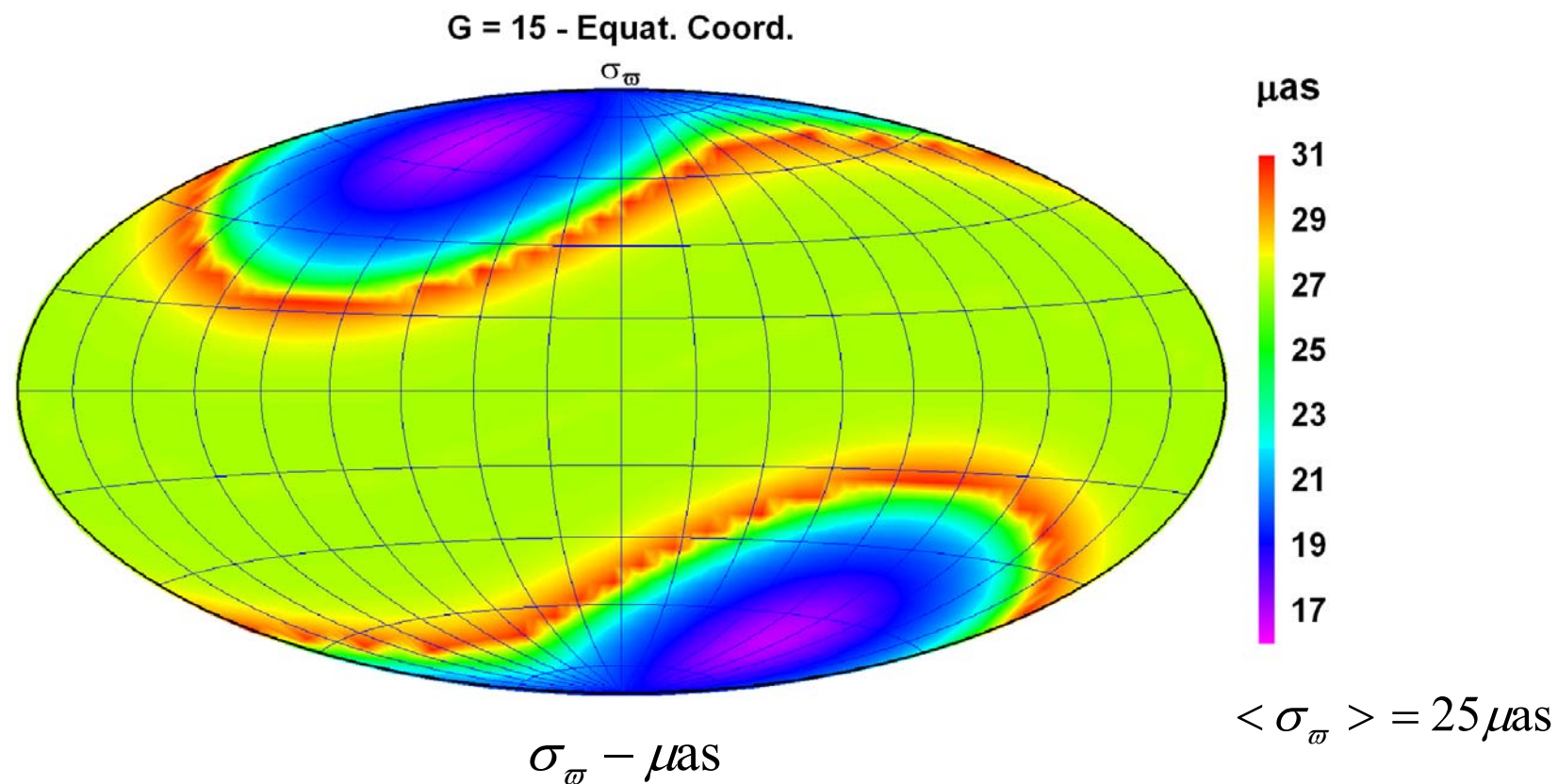
$$\langle \sigma_{\alpha}^* \rangle = 21 \mu\text{as}$$

$$\langle \sigma_{\delta} \rangle = 18 \mu\text{as}$$



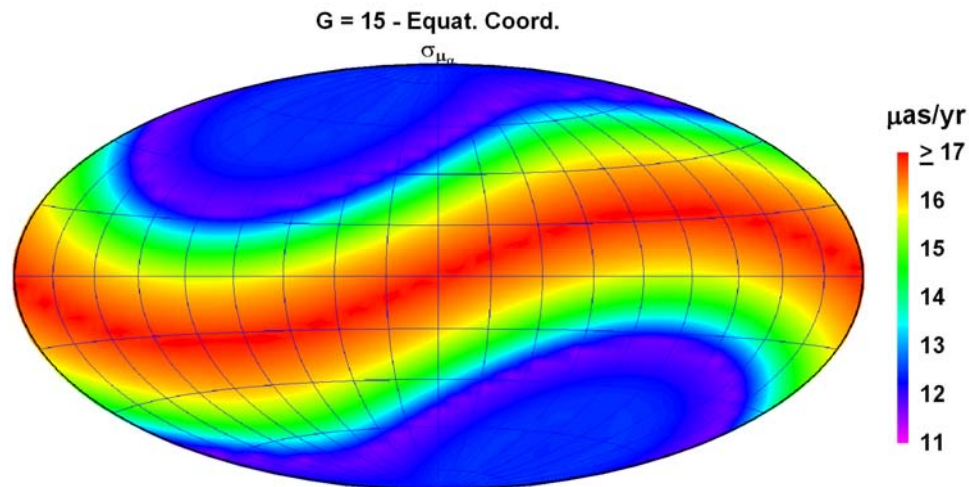
Equatorial coordinates

- Plot for $G = 15$, but scalable to other magnitudes



Equatorial coordinates

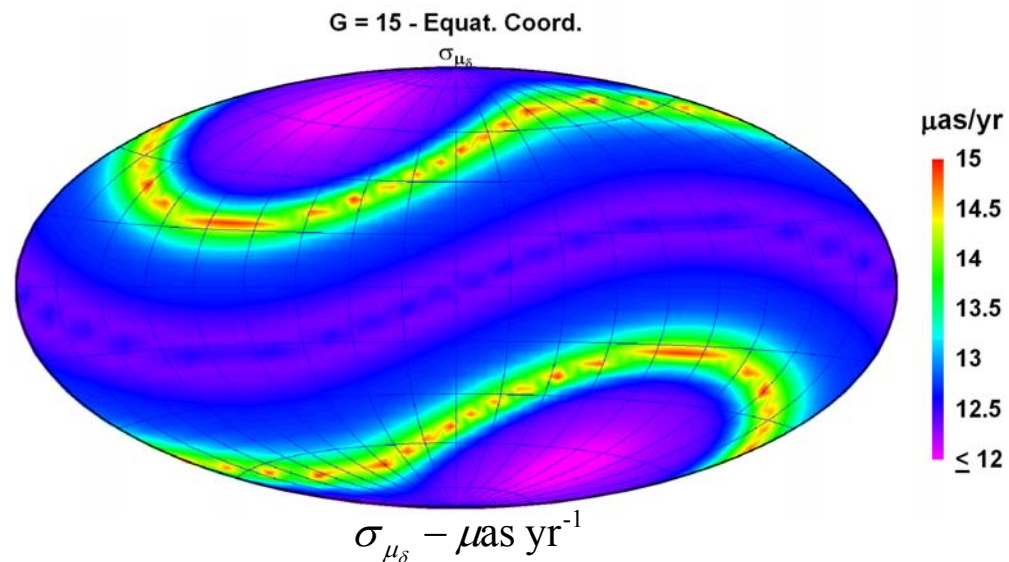
- Plots for $G = 15$, but scalable to other magnitudes



$$\langle \sigma_{\mu_\alpha}^* \rangle = 15 \mu\text{as yr}^{-1}$$

$$\sigma_{\mu_\alpha}^* - \mu\text{as yr}^{-1}$$

$$\langle \sigma_{\mu_\delta} \rangle = 13 \mu\text{as yr}^{-1}$$

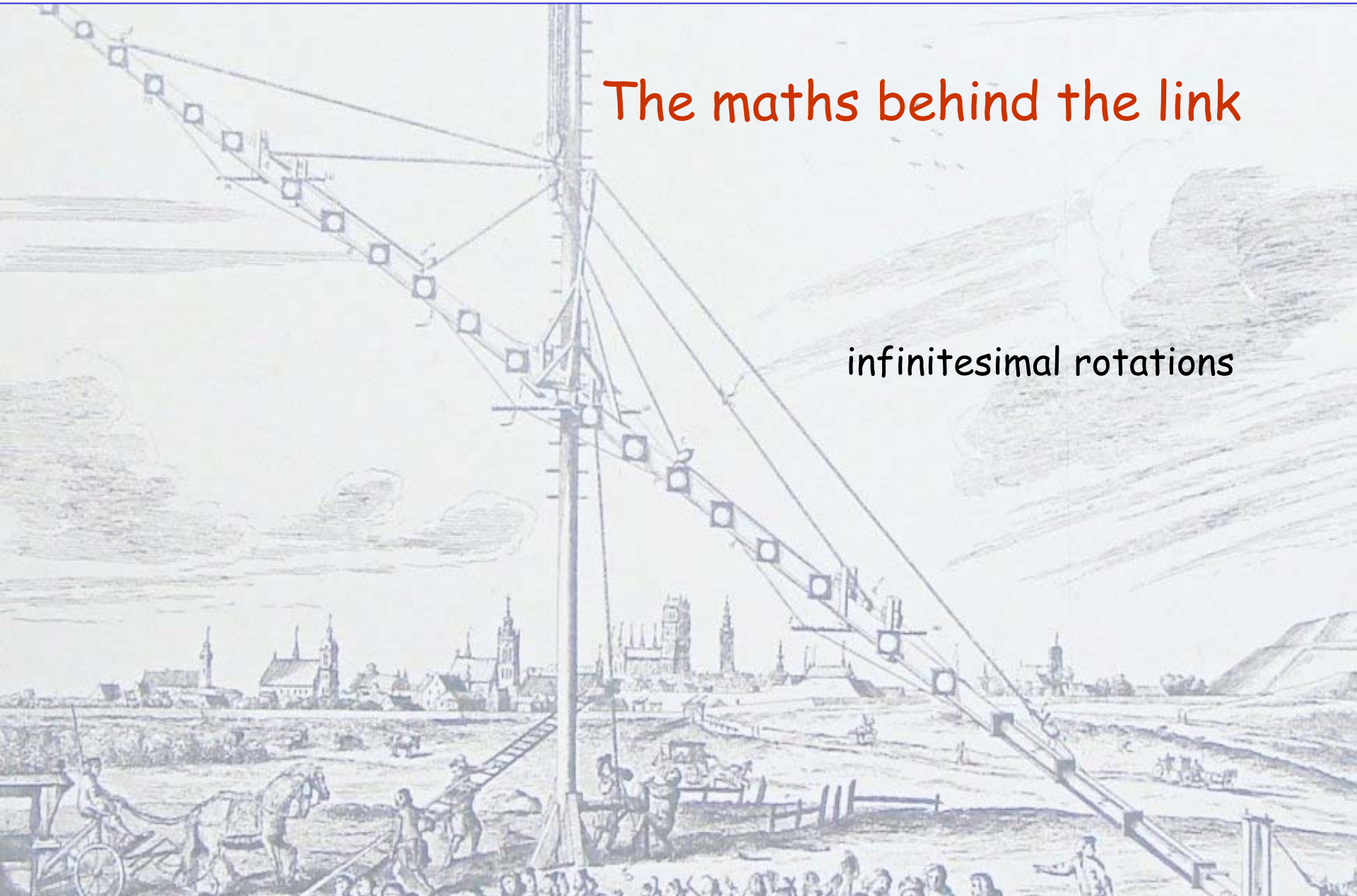


$$\sigma_{\mu_\delta} - \mu\text{as yr}^{-1}$$

Equatorial coordinates

The maths behind the link

infinitesimal rotations



- Local frame

\mathbf{u} unit vector in direction \mathbf{OM}

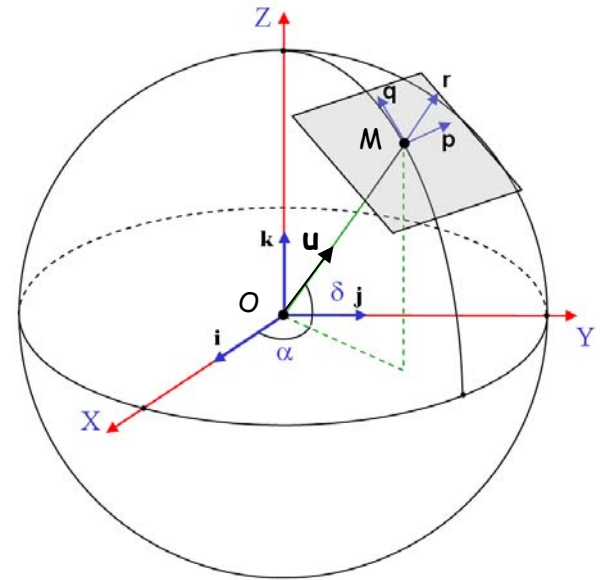
$$\mathbf{u} = \begin{bmatrix} \cos \alpha \cos \delta \\ \sin \alpha \cos \delta \\ \sin \delta \end{bmatrix}$$

$$\mathbf{p} = \frac{1}{\cos \delta} \frac{d\mathbf{u}}{d\alpha} = \begin{bmatrix} -\sin \alpha \\ \cos \alpha \\ 0 \end{bmatrix} \quad \mathbf{q} = \frac{d\mathbf{u}}{d\delta} = \begin{bmatrix} -\cos \alpha \sin \delta \\ -\sin \alpha \sin \delta \\ \cos \delta \end{bmatrix}$$

$[\mathbf{p}, \mathbf{q}, \mathbf{r}]$ direct orthonormal triad

$$[\mathbf{i}, \mathbf{j}, \mathbf{k}] \Rightarrow [\mathbf{r}, \mathbf{p}, \mathbf{q}] : R_2(-\delta)R_3(\alpha)$$

$$\mathbf{V}_{[\mathbf{r}, \mathbf{p}, \mathbf{q}]} = \begin{bmatrix} \cos \alpha \cos \delta & \sin \alpha \cos \delta & \sin \delta \\ -\sin \alpha & \cos \alpha & 0 \\ -\cos \alpha \sin \delta & -\sin \alpha \sin \delta & \cos \delta \end{bmatrix} \mathbf{V}_{[\mathbf{i}, \mathbf{j}, \mathbf{k}]}$$



- Rotation ω $|\omega| \ll 1$

$$\boxed{d\mathbf{u} = \omega \times \mathbf{r}} \quad (\text{active form})$$

fixed frame

$$d\mathbf{u} = \begin{bmatrix} 0 & -\omega_z & \omega_y \\ \omega_z & 0 & -\omega_x \\ -\omega_y & \omega_x & 0 \end{bmatrix} \begin{matrix} x \\ y \\ z \end{matrix} \quad (\text{passive form})$$

local frame

$$d\mathbf{u} \cdot \mathbf{p} = (\omega \times \mathbf{r}) \cdot \mathbf{p} = (\mathbf{r} \times \mathbf{p}) \cdot \omega = \omega \cdot \mathbf{q}$$

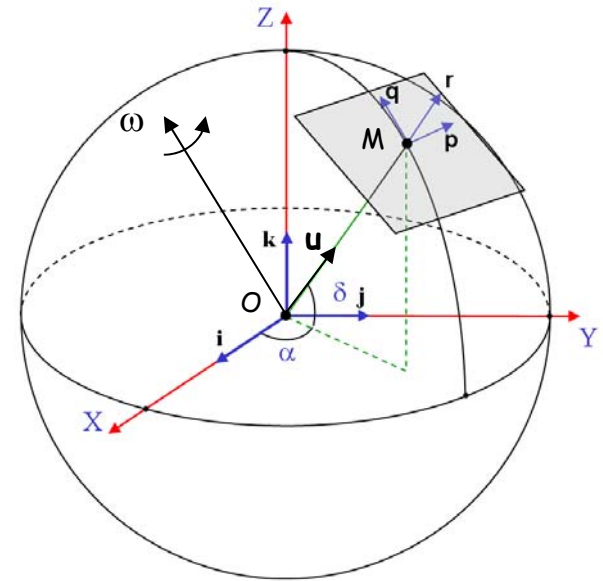
$$d\mathbf{u} \cdot \mathbf{q} = (\omega \times \mathbf{r}) \cdot \mathbf{q} = (\mathbf{r} \times \mathbf{q}) \cdot \omega = -\omega \cdot \mathbf{p}$$

$$\omega \cdot \mathbf{q} = \omega \cdot \frac{d\mathbf{u}}{d\delta} = \frac{d(\omega \cdot \mathbf{u})}{d\delta} \quad \omega \cdot \mathbf{p} = \omega \cdot \frac{1}{\cos \delta} \frac{d\mathbf{u}}{d\alpha} = \frac{1}{\cos \delta} \frac{d(\omega \cdot \mathbf{u})}{d\alpha}$$

$$\omega \cdot \mathbf{u} = \omega_x u_x + \omega_y u_y + \omega_z u_z = \omega_x \cos \alpha \cos \delta + \omega_y \sin \alpha \cos \delta + \omega_z \sin \delta$$

$$d\mathbf{u} \cdot \mathbf{p} = -\omega_x \cos \alpha \sin \delta - \omega_y \sin \alpha \sin \delta + \omega_z \cos \delta$$

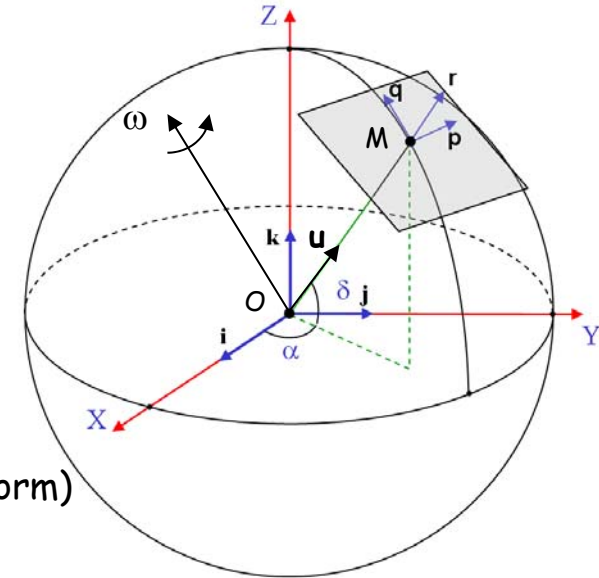
$$d\mathbf{u} \cdot \mathbf{q} = \omega_x \sin \alpha - \omega_y \cos \alpha$$



$$d\mathbf{u} \cdot \mathbf{p} = -\omega_x \cos \alpha \sin \delta - \omega_y \sin \alpha \sin \delta + \omega_z \cos \delta$$

$$d\mathbf{u} \cdot \mathbf{q} = \omega_x \sin \alpha - \omega_y \cos \alpha$$

$$\begin{bmatrix} \Delta\alpha \cos \delta \\ \Delta\delta \end{bmatrix} = \begin{bmatrix} \cos \alpha \sin \delta & \sin \alpha \sin \delta & -\cos \delta \\ -\sin \alpha & \cos \alpha & 0 \end{bmatrix} \begin{bmatrix} \omega_x \\ \omega_y \\ \omega_z \end{bmatrix} \quad (\text{passive form})$$



- Relationship between two frames related by a small, **static** rotation

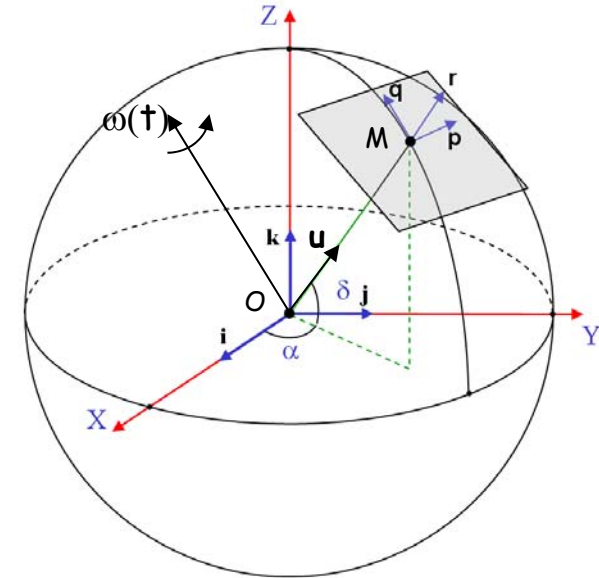
- correspondence between two catalogues of the **same** sources given in each frame
- one can use the matrix to rotate one catalogue from one frame to the other
 - ω is known
- one can use the matrix as condition equations to derive ω
 - the $(\Delta\alpha \cos\delta, \Delta\delta)_i$ are known

- Non static form

- rotation $\omega(t)$ $\omega(t) = \boldsymbol{\varepsilon} + \boldsymbol{\Omega}(t - t_0)$

$$\boldsymbol{\varepsilon}: [\varepsilon_x, \varepsilon_y, \varepsilon_z] \quad ; \quad \boldsymbol{\Omega}: [\Omega_x, \Omega_y, \Omega_z]$$

$$\begin{matrix} \Delta\alpha^*(t) \\ \Delta\delta(t) \end{matrix} = \begin{bmatrix} \cos\alpha \sin\delta & \sin\alpha \sin\delta & -\cos\delta \\ -\sin\alpha & \cos\alpha & 0 \end{bmatrix} \begin{bmatrix} \omega_x(t) \\ \omega_y(t) \\ \omega_z(t) \end{bmatrix}$$



with two catalogue or two frames

$$\Delta\delta(t) = \delta_2(t) - \delta_1(t) = [\delta_0 + \mu_\delta(t - t_0)]_2 - [\delta_0 + \mu_\delta(t - t_0)]_1 \quad \boxed{= \Delta\delta_0 + \Delta\mu_\delta(t - t_0)}$$

similarly:

$$\boxed{\Delta\alpha^*(t) = \Delta\alpha^*_0 + \Delta\mu_\alpha(t - t_0)}$$

$$\boldsymbol{\omega}(t) = \boldsymbol{\varepsilon} + \boldsymbol{\Omega}(t - t_0)$$

the transformation reads:

$$\begin{array}{l} \Delta\alpha^*(t) \\ \Delta\delta(t) \end{array} = \begin{bmatrix} \cos\alpha \sin\delta & \sin\alpha \sin\delta & -\cos\delta \\ -\sin\alpha & \cos\alpha & 0 \end{bmatrix} \begin{bmatrix} \varepsilon_x + \Omega_x(t - t_0) \\ \varepsilon_y + \Omega_y(t - t_0) \\ \varepsilon_z + \Omega_z(t - t_0) \end{bmatrix}$$

and by identification:

$$\begin{array}{l} \Delta\alpha_0^* \\ \Delta\delta_0 \end{array} = \begin{bmatrix} \cos\alpha \sin\delta & \sin\alpha \sin\delta & -\cos\delta \\ -\sin\alpha & \cos\alpha & 0 \end{bmatrix} \begin{bmatrix} \varepsilon_x \\ \varepsilon_y \\ \varepsilon_z \end{bmatrix}$$

$$\begin{array}{l} \Delta\mu_\alpha \\ \Delta\mu_\delta \end{array} = \begin{bmatrix} \cos\alpha \sin\delta & \sin\alpha \sin\delta & -\cos\delta \\ -\sin\alpha & \cos\alpha & 0 \end{bmatrix} \begin{bmatrix} \Omega_x \\ \Omega_y \\ \Omega_z \end{bmatrix}$$

- Two sets of coordinates of the same sources
 - one from Gaia
 - one from a reference catalogue
 - ICRF2 for the orientation, QSO catalogue for the spin
 - if the differences can be approximated by a time dependant rotation, for each source one has,

$$\Delta\alpha_{\alpha}^* \approx \varepsilon_x \cos \alpha \sin \delta + \varepsilon_y \sin \alpha \sin \delta - \varepsilon_x \cos \delta$$

$$\Delta\delta_{\delta} \approx -\varepsilon_x \sin \alpha + \varepsilon_y \cos \alpha$$

$$\Delta\mu_{\alpha}^* \approx \Omega_x \cos \alpha \sin \delta + \Omega_y \sin \alpha \sin \delta - \Omega_x \cos \delta$$

$$\Delta\mu_{\delta} \approx -\Omega_x \sin \alpha + \Omega_y \cos \alpha$$

- Solution for Ω and ε with least-squares

- The normal matrix depends only on the source distribution and accuracy (\sim magnitude)

$$N_{xx} = \sum W_{\alpha} \cos^2 \alpha \sin^2 \delta + \sum W_{\delta} \sin^2 \alpha$$

$$N_{xy} = \sum W_{\alpha} \sin \alpha \cos \alpha \sin^2 \delta - \sum W_{\delta} \sin \alpha \cos \alpha$$

$$N_{yy} = \sum W_{\alpha} \sin^2 \alpha \sin^2 \delta + \sum W_{\delta} \cos^2 \alpha$$

$$N_{yz} = -\sum W_{\alpha} \sin \alpha \cos \delta \sin \delta$$

$$N_{zz} = \sum W_{\alpha} \cos^2 \delta$$

$$N_{zx} = -\sum W_{\alpha} \cos \alpha \cos \delta \sin \delta$$

- With a relatively uniform distribution in each bin of magnitude

$$\langle \cos^2 \alpha \rangle \approx \langle \sin^2 \alpha \rangle \approx 1/2$$

$$\langle \sin \alpha \rangle \approx \langle \cos \alpha \rangle \approx \langle \sin \alpha \cos \alpha \rangle \approx 0$$

$$\langle \sin^2 \delta \rangle \approx 1/3 \quad \langle \cos^2 \delta \rangle \approx 2/3$$

- the normal matrix is nearly diagonal

$$\text{cov}(\boldsymbol{\epsilon}) = \mathbf{N}^{-1}$$

alignment sources

$$\text{cov}(\boldsymbol{\Omega}) = \mathbf{N}^{-1}$$

spin sources

- Light and radio emission centers could be different
 - there is at the moment no real evidence below ~ 10 mas level
 - physically with QSO models one has good reasons to assume there is an offset
 - how big is an open question

model without offset

$$\begin{matrix} \Delta\alpha_0^* \\ \Delta\delta_0 \end{matrix} = \begin{bmatrix} \cos\alpha \sin\delta & \sin\alpha \sin\delta & -\cos\delta \\ -\sin\alpha & \cos\alpha & 0 \end{bmatrix} \begin{bmatrix} \varepsilon_x \\ \varepsilon_y \\ \varepsilon_z \end{bmatrix} + \begin{bmatrix} N_\alpha \\ N_\delta \end{bmatrix}$$

noise

model with offset

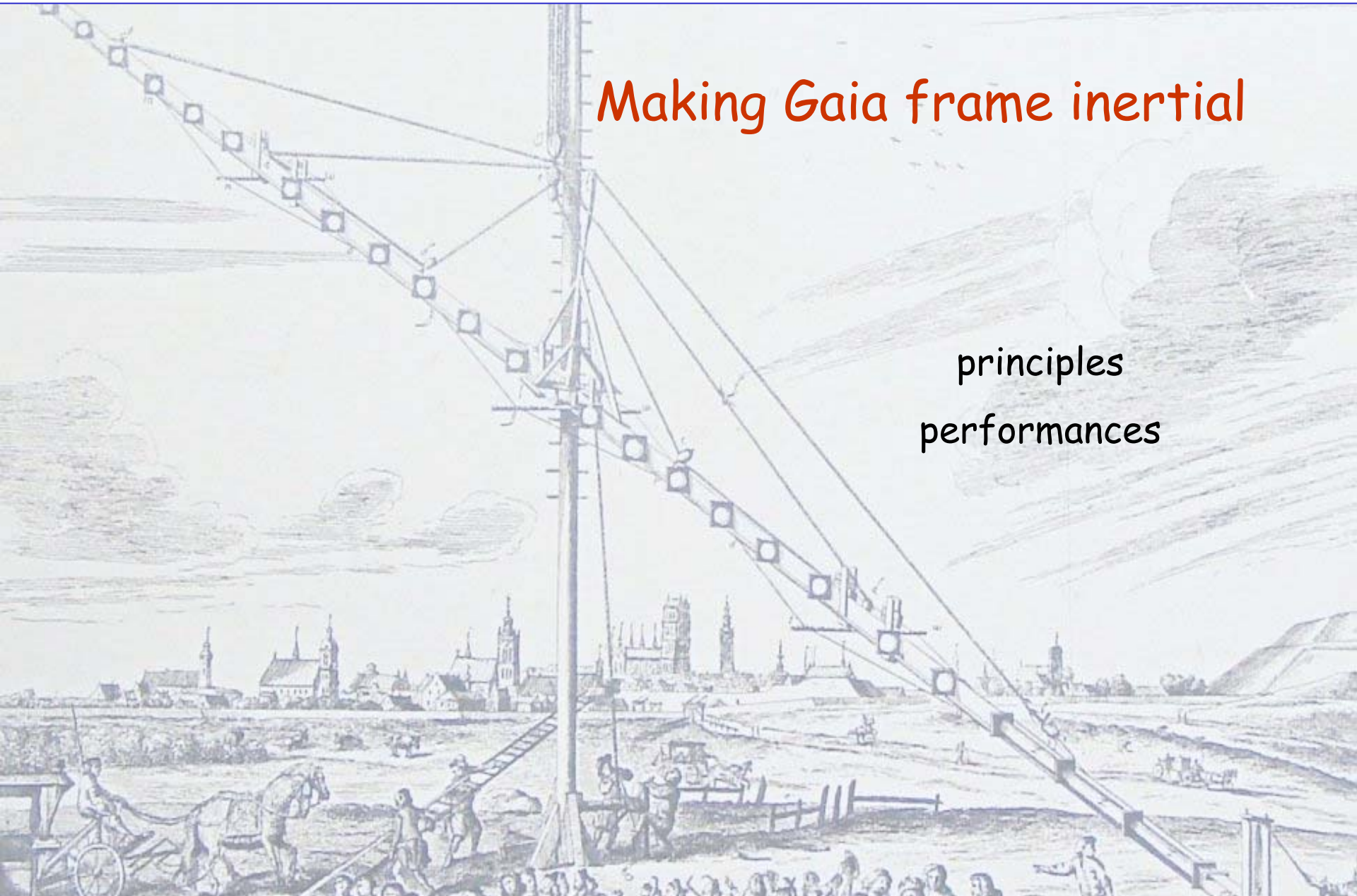
$$\begin{matrix} \Delta\alpha_0^* \\ \Delta\delta_0 \end{matrix} = \begin{bmatrix} \cos\alpha \sin\delta & \sin\alpha \sin\delta & -\cos\delta \\ -\sin\alpha & \cos\alpha & 0 \end{bmatrix} \begin{bmatrix} \varepsilon_x \\ \varepsilon_y \\ \varepsilon_z \end{bmatrix} + \begin{bmatrix} O_\alpha \\ O_\delta \end{bmatrix} + \begin{bmatrix} N_\alpha \\ N_\delta \end{bmatrix}$$

offset noise

- The offset vector \mathbf{O} is probably of random nature from source to source

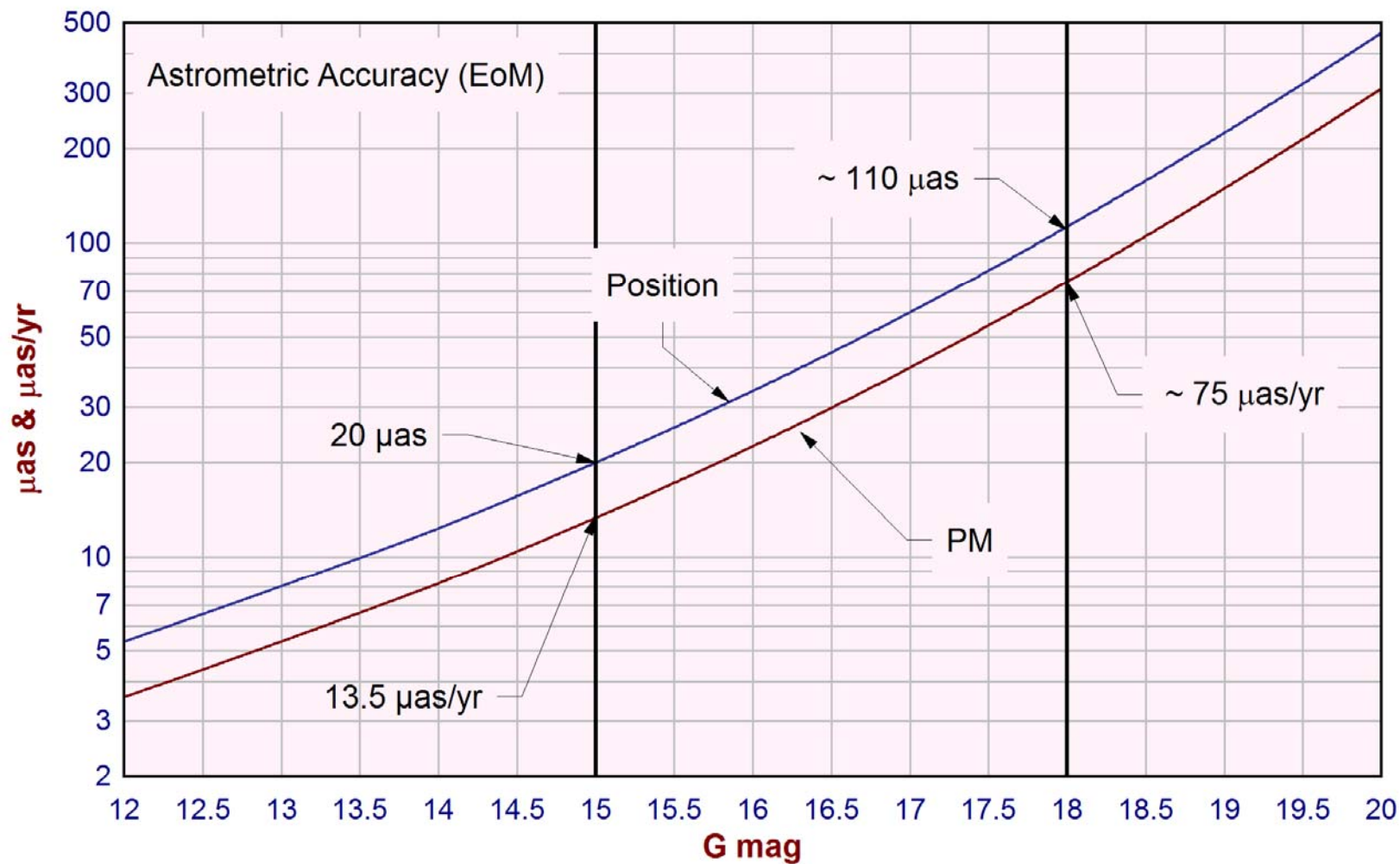
Making Gaia frame inertial

principles
performances

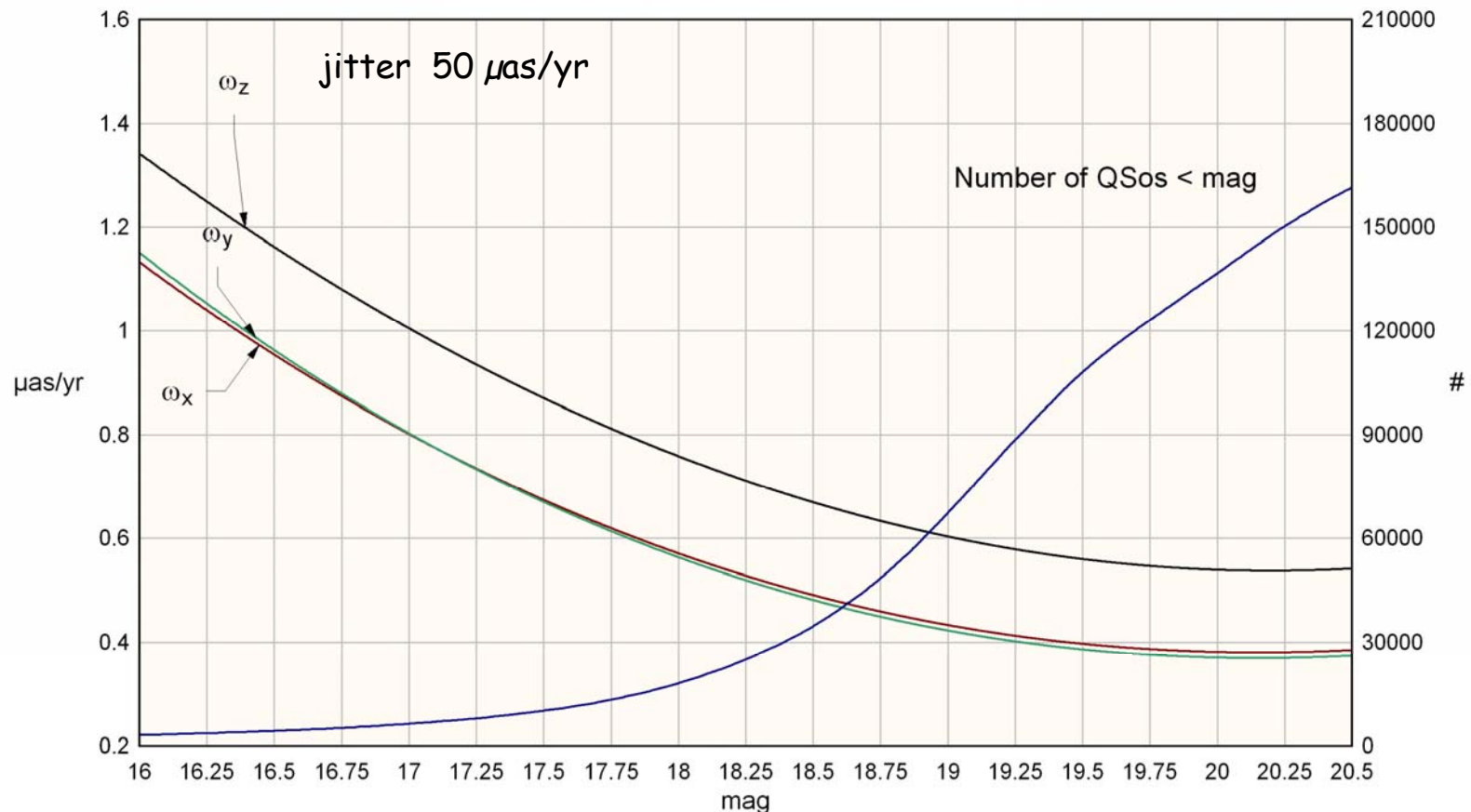


- **Assumption**
 - the only source of transverse motion in the Gaia solution comes from the free spin
- **Material**
 - LQAC known sources (150,000 sources $G < 20$)
 - Gaia simulated QSO catalogue (550,000 sources $G < 20$)
- **Model fitting**
 - A global spin ω ($\omega_x \ \omega_y \ \omega_z$) on the QSO proper motions of the Gaia unrotated frame
 - the covariance matrix is computed for each bin of magnitude
- **Application**
 - all proper motion (stars, QSOs) are corrected for this spin pattern

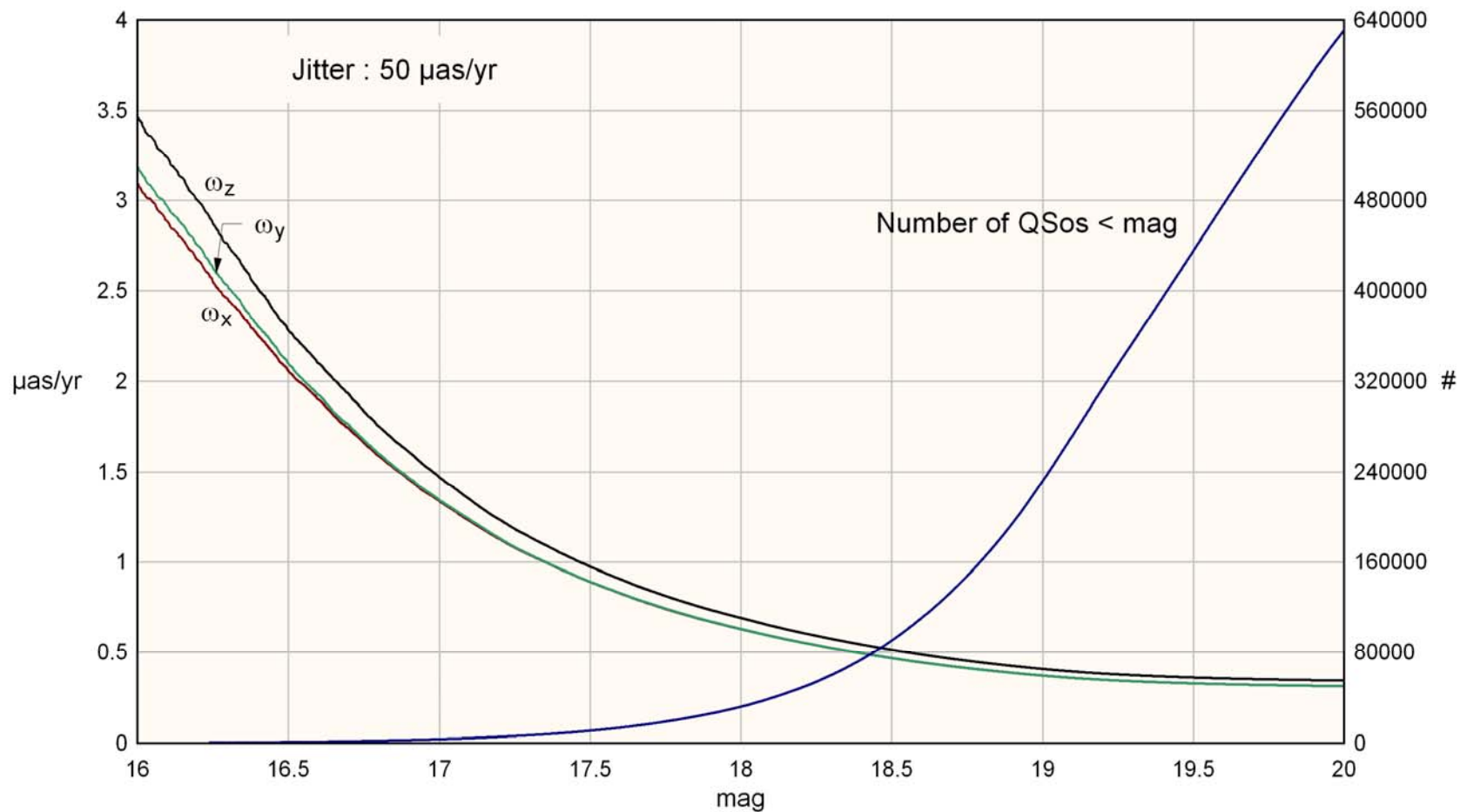
- So far no systematic transverse motion detected
 - QSOs have fixed comoving coordinates
- If $V_{\dagger} \sim H_0 D \rightarrow \mu \sim 10 \mu\text{as/yr}$
 - VLBI in 20 yrs with $\sigma_{\text{pos}} \sim 1 \text{ mas} \rightarrow \mu < 50 \mu\text{as}$
 - but sub-mas structure instabilities (P. Charlot, 2003)
- Other sources :
 - microlensing $P = 10^{-6}$ (Belokurov) \rightarrow only a handful
 - matter ejection, superluminous motion
 - Variable galactic aberration
 - Macrolensing $P = 10^{-2}$ (Mignard, 2003) \rightarrow long timescale
 - Accelerated motion in the local group
 - binary QSOs ?



- Spin covariance matrix computed when QSOs are constrained to have no overall motion
- The plot shows the standard error in ω



- Spin covariance matrix computed when QSOs are constrained to have no overall motion
- The plot shows the standard error in ω



We aim at a final result on the reference frame and the acceleration at $0.3 \mu\text{as}/\text{yr}$ level:

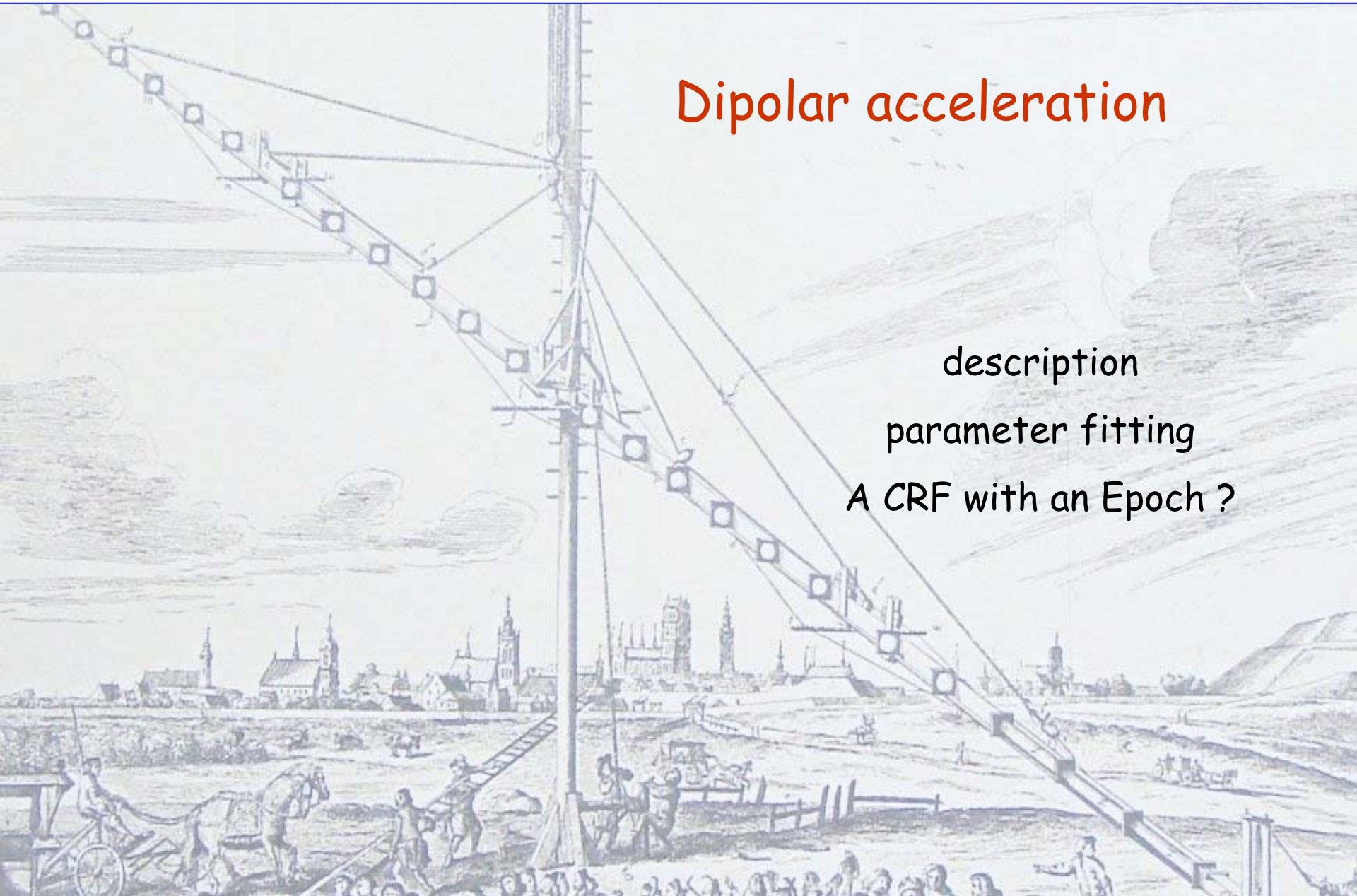
This is 1/1000 of the astrometric accuracy of the faintest sources

Dipolar acceleration

description

parameter fitting

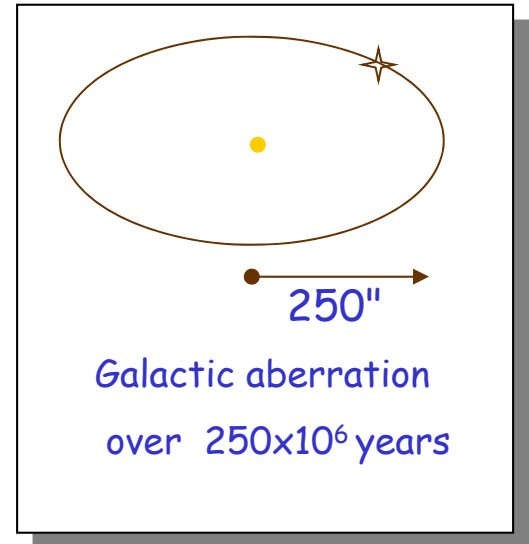
A CRF with an Epoch ?



- The solar system is in motion in the Galaxy, $V \sim 220 \text{ km s}^{-1}$
 - constant aberration of $\sim 250''$ for the QSO wrt to comoving frame
 - not detectable (principle of relativity)
 - $\delta \mathbf{u} = \mathbf{v}/c$

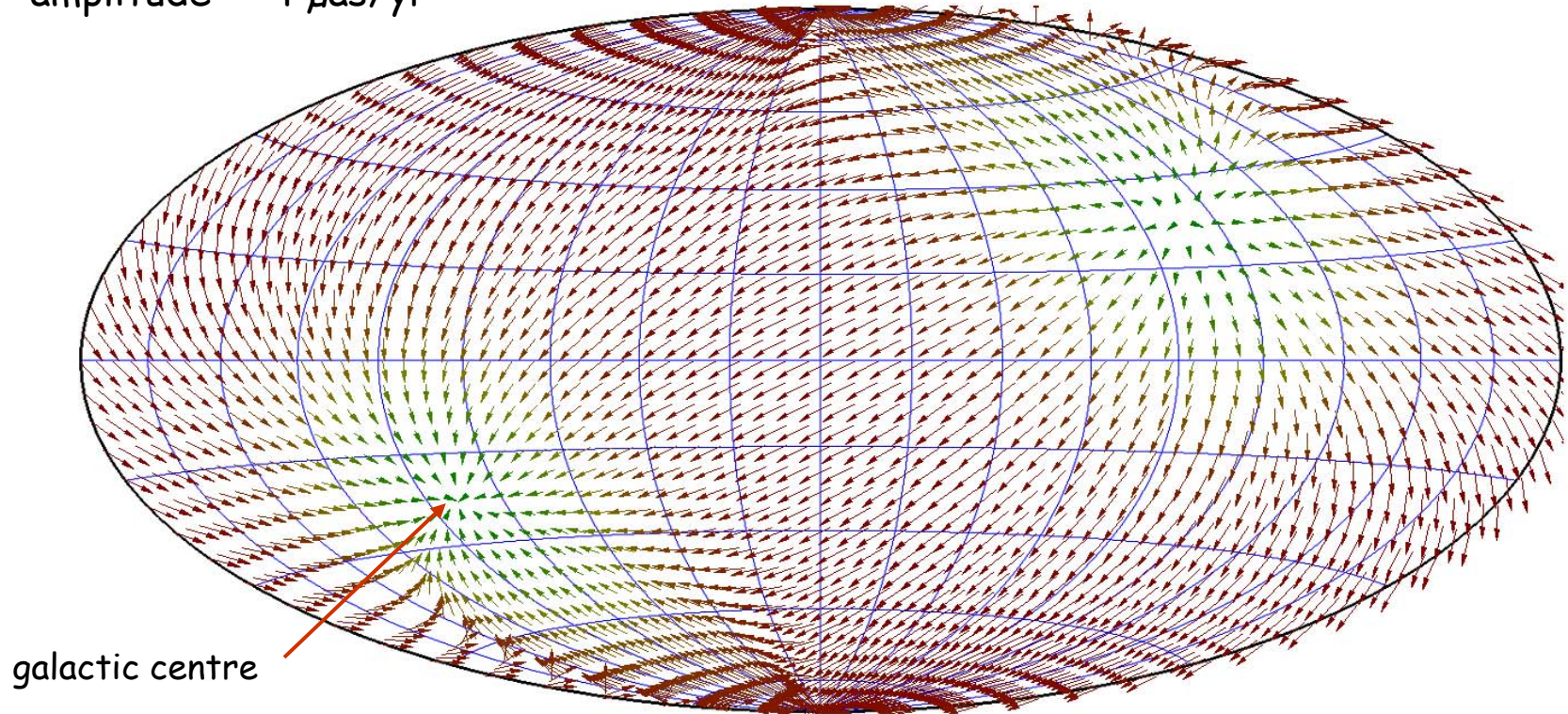
- But the solar motion is not uniform

- \sim circular motion
 - radius $R \sim 8.5 \text{ kpc}$ and period $250 \times 10^6 \text{ yrs}$
- the aberration is then variable
 - one sees a very small arc on the aberration ellipse



$$\delta \mu = \frac{d(\delta \mathbf{u})}{dt} = \frac{\mathbf{\Gamma}}{c} = \frac{V^2}{cR} \approx 4 \mu\text{as} / \text{yr}$$

- Plot in equatorial coordinates
 - amplitude $\sim 4 \mu\text{as/yr}$



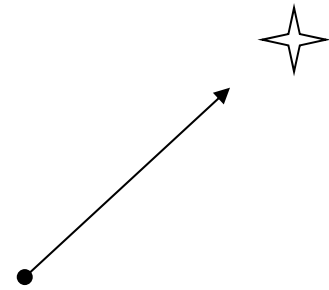
$$\mu_{\alpha} \cos \delta = -\frac{\Gamma_x}{c} \sin \alpha + \frac{\Gamma_y}{c} \cos \alpha$$

$$\mu_{\delta} = -\frac{\Gamma_x}{c} \sin \delta \cos \alpha - \frac{\Gamma_y}{c} \sin \delta \sin \alpha + \frac{\Gamma_z}{c} \cos \delta$$



- For any acceleration of the SS wrt Quasars :

$$\frac{d\mathbf{u}}{dt} = \frac{\mathbf{\Gamma}}{c} - \left(\frac{\mathbf{\Gamma}}{c} \cdot \mathbf{u}\right)\mathbf{u}$$



Observations

$$\mu_{\alpha} \cos \delta = -\frac{\Gamma_x}{c} \sin \alpha + \frac{\Gamma_y}{c} \cos \alpha$$

$$\mu_{\delta} = -\frac{\Gamma_x}{c} \sin \delta \cos \alpha - \frac{\Gamma_y}{c} \sin \delta \sin \alpha + \frac{\Gamma_z}{c} \cos \delta$$

- Equations similar to global rotation.
- Precision of $\sim 0.3 \mu\text{as/yr}$ (2 prad/yr) on Γ/c
 $= 0.2 \times 10^{-10} \text{ m s}^{-2}$ ($\gamma \text{ Pionner}/40$)
- Galactic rotation ($\mu \sim 4 \mu\text{as/yr}$)
- Acceleration of the Local Group \rightarrow CDM ?

- The two fields are **globally** orthogonal on the sphere

$$\boldsymbol{\mu}_a = \mathbf{u} \times (\mathbf{a} \times \mathbf{u}) = \mathbf{a} - (\mathbf{a} \cdot \mathbf{u})\mathbf{u}$$

$$\boldsymbol{\mu}_\omega = \boldsymbol{\Omega} \times \mathbf{u}$$

$$\boldsymbol{\mu}_a \cdot \boldsymbol{\mu}_\omega = (\boldsymbol{\Omega} \times \mathbf{u}) \cdot (\mathbf{a} - (\mathbf{a} \cdot \mathbf{u})\mathbf{u}) = (\boldsymbol{\Omega} \times \mathbf{u}) \cdot \mathbf{a} = (\mathbf{a} \times \boldsymbol{\Omega}) \cdot \mathbf{u}$$

- If \mathbf{a} and $\boldsymbol{\Omega}$ are parallel, then local orthogonality
- Otherwise

$$\int_S \boldsymbol{\mu}_a \cdot \boldsymbol{\mu}_\omega dS = \int_S (\mathbf{a} \times \boldsymbol{\Omega}) \cdot \mathbf{u} dS = (\mathbf{a} \times \boldsymbol{\Omega}) \cdot \int_S \mathbf{u} dS \equiv 0$$

- Orthogonal only on the average on the sphere
 - therefore, one must solve **simultaneously** for both vectors

- The galactic acceleration entails a systematic transverse motion of the QSOs

$$\frac{d\mathbf{u}}{dt} = \boldsymbol{\mu} = \mathbf{u} \times \left(\frac{\boldsymbol{\Gamma}}{c} \times \mathbf{u} \right) = \frac{\boldsymbol{\Gamma}}{c} - \left(\frac{\boldsymbol{\Gamma}}{c} \cdot \mathbf{u} \right) \mathbf{u}$$

$$\boldsymbol{\mu} \cdot \mathbf{p} = \frac{\boldsymbol{\Gamma} \cdot \mathbf{p}}{c} \quad \boldsymbol{\mu} \cdot \mathbf{q} = \frac{\boldsymbol{\Gamma} \cdot \mathbf{q}}{c} \quad \text{since} \quad \mathbf{u} \cdot \mathbf{p} = \mathbf{u} \cdot \mathbf{q} \equiv 0$$

with $\mathbf{a} = \frac{\boldsymbol{\Gamma}}{c}$:

$$\begin{matrix} \Delta\mu_a^* \\ \Delta\mu_\delta \end{matrix} = \begin{bmatrix} \cos\alpha \sin\delta & \sin\alpha \sin\delta & -\cos\delta \\ -\sin\alpha & \cos\alpha & 0 \end{bmatrix} \begin{bmatrix} \Omega_x \\ \Omega_y \\ \Omega_z \end{bmatrix} + \begin{bmatrix} -\sin\alpha & \cos\alpha & 0 \\ -\cos\alpha \sin\delta & -\sin\alpha \sin\delta & \cos\delta \end{bmatrix} \begin{bmatrix} a_x \\ a_y \\ a_z \end{bmatrix} + \begin{bmatrix} N_{\mu_a} \\ N_{\mu_\delta} \end{bmatrix}$$

noise



- Given the near perfect orthogonality the normal matrix is block diagonal

$$\begin{bmatrix} \mathbf{N} & \mathbf{0} \\ \mathbf{0} & \mathbf{N} \end{bmatrix} \quad \text{and} \quad \text{cov}(\boldsymbol{\Omega}) \sim \text{cov}(\mathbf{a}) = \mathbf{N}^{-1}$$

- Proper motions seen as a vector field on S_2
- Applicable to stars and QSOs
- Expansion in Vector Spherical Harmonics \mathbf{T}_{lm} , \mathbf{S}_{lm}

$$\mathbf{V}(\alpha, \delta) = V_\alpha \mathbf{e}_\alpha + V_\delta \mathbf{e}_\delta = \sum_{l=1}^{l=L} \sum_{m=-l}^{m=l} (t_{lm} \mathbf{T}_{lm} + s_{lm} \mathbf{S}_{lm})$$

$l = 1 - \mathbf{S}_{1m}$	Global rotation
$l = 1 - \mathbf{T}_{1m}$	Solar system acceleration
$l > 1 - \mathbf{S}_{1m} \text{ \& } \mathbf{T}_{1m}$	Stochastic field of GW

Will the Gaia-CRF have an Epoch ?

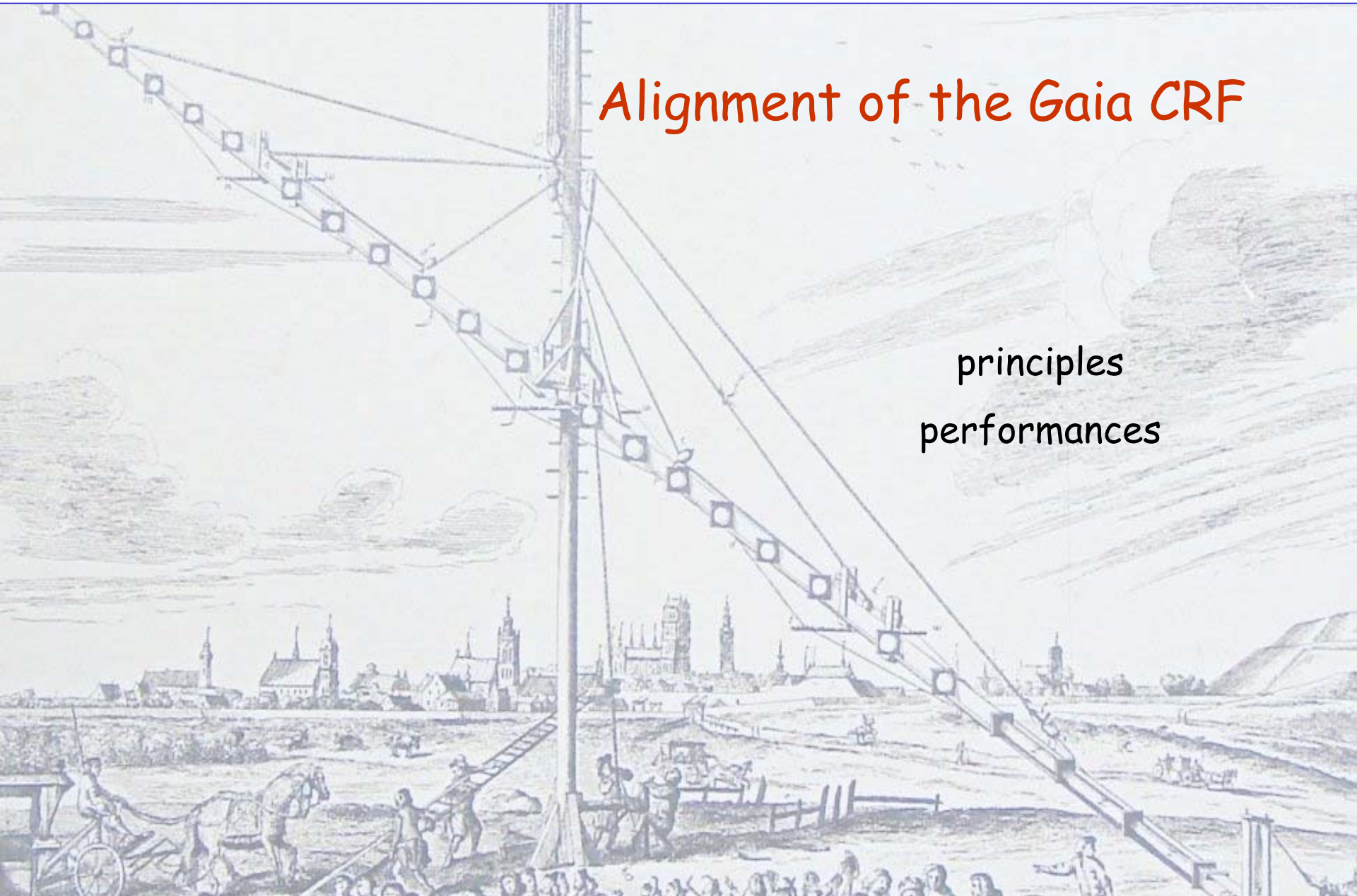


- The radio ICRF is not associated to an epoch
 - defining QSOs have fixed celestial coordinates
 - they are not epoch dependant
 - the define axis direction 'for ever'
 - time is no involved in the process
- A stellar reference frame is defined at a particular epoch
 - defining stars coordinates come with their proper motion
 - the PMs are part of the fundamental catalogue
 - each star comes with a particular PM and its uncertainty
 - with N stars, there are $2N$ parameters needed
 - the system degrades due to the limited uncertainty of the PMs
 - accuracy in position and annual PM are similar

- What about the Gaia-CRF
 - QSOs have a systematic proper motion of ~ 4 mas/yr
 - But these are not individual PM, but the result of a systematic pattern
 - only 3 parameters are required to maintain the system
 - the accuracy should be < 0.5 mas/yr
 - Individual positions of the primary sources will have an accuracy of ~ 80 mas
 - degradation will be very very slow
- Therefore : the Gaia-CRF will have an epoch attached to it
 - but it has very different meaning as for a stellar reference frame
- How to avoid it: take the origin at the galactic centre !
 - this is for the future

Alignment of the Gaia CRF

principles
performances

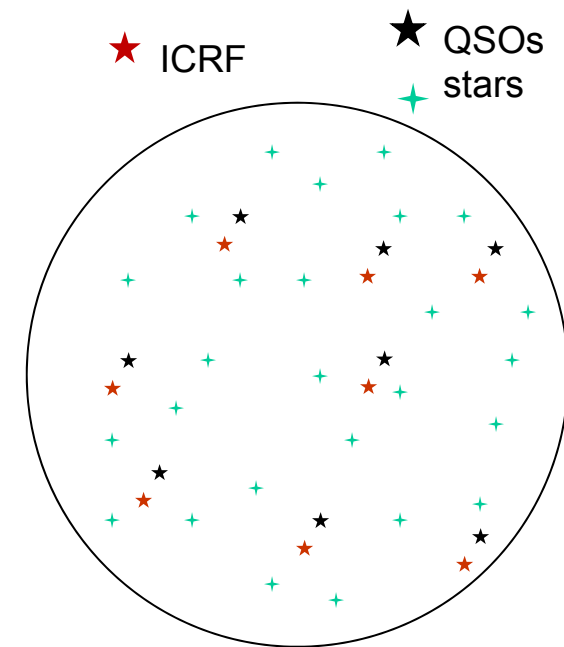


- Orientation is performed by minimizing the distances between Gaia positions and ICRF positions of common sources
 - GCRF needs to be aligned to ICRF
 - we have one infinitesimal rotations to fit($\varepsilon_x, \varepsilon_y, \varepsilon_z$)

- ICRF sources are observed by Gaia

- $\sim 1500 G < 20$, $200 G < 18$ - $\sigma_{\text{Gaia}} < 100 \mu\text{as}$
- GCRF can be aligned to QSOs by a rotation
- accuracy

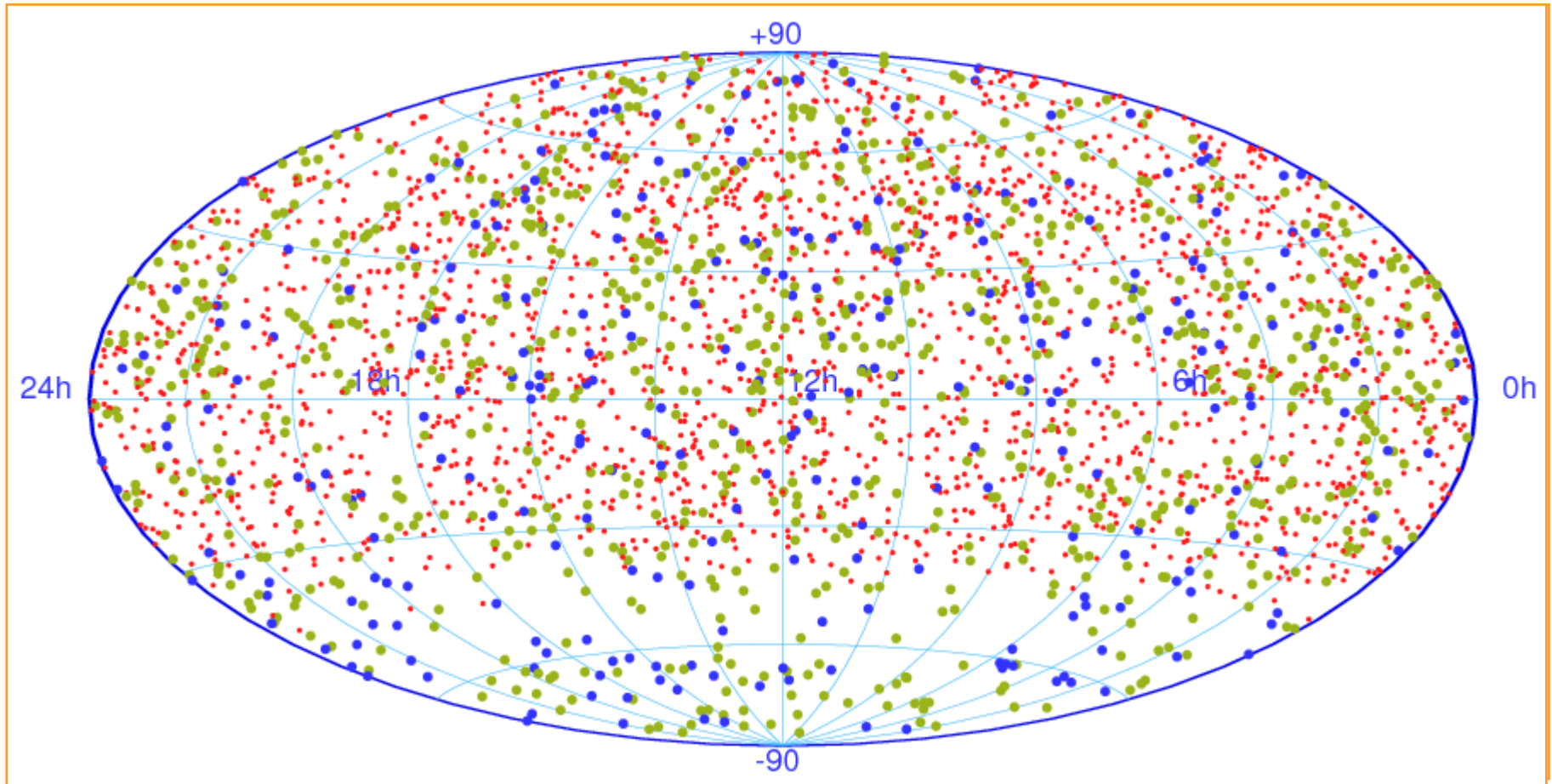
$$\sigma_{\text{align}} \approx \frac{\sqrt{\sigma_{\text{Gaia}}^2 + \sigma_{\text{ICRF}}^2}}{\sqrt{N_{\text{QSO}}}} < 10 \mu\text{as}$$



$\sigma \sim 50$ to $150 \mu\text{as}$

$\sigma \sim 0.2$ to 2 mas

$\sigma \sim 0.5$ to 10 mas



● defining (294)

● VLBI (923)

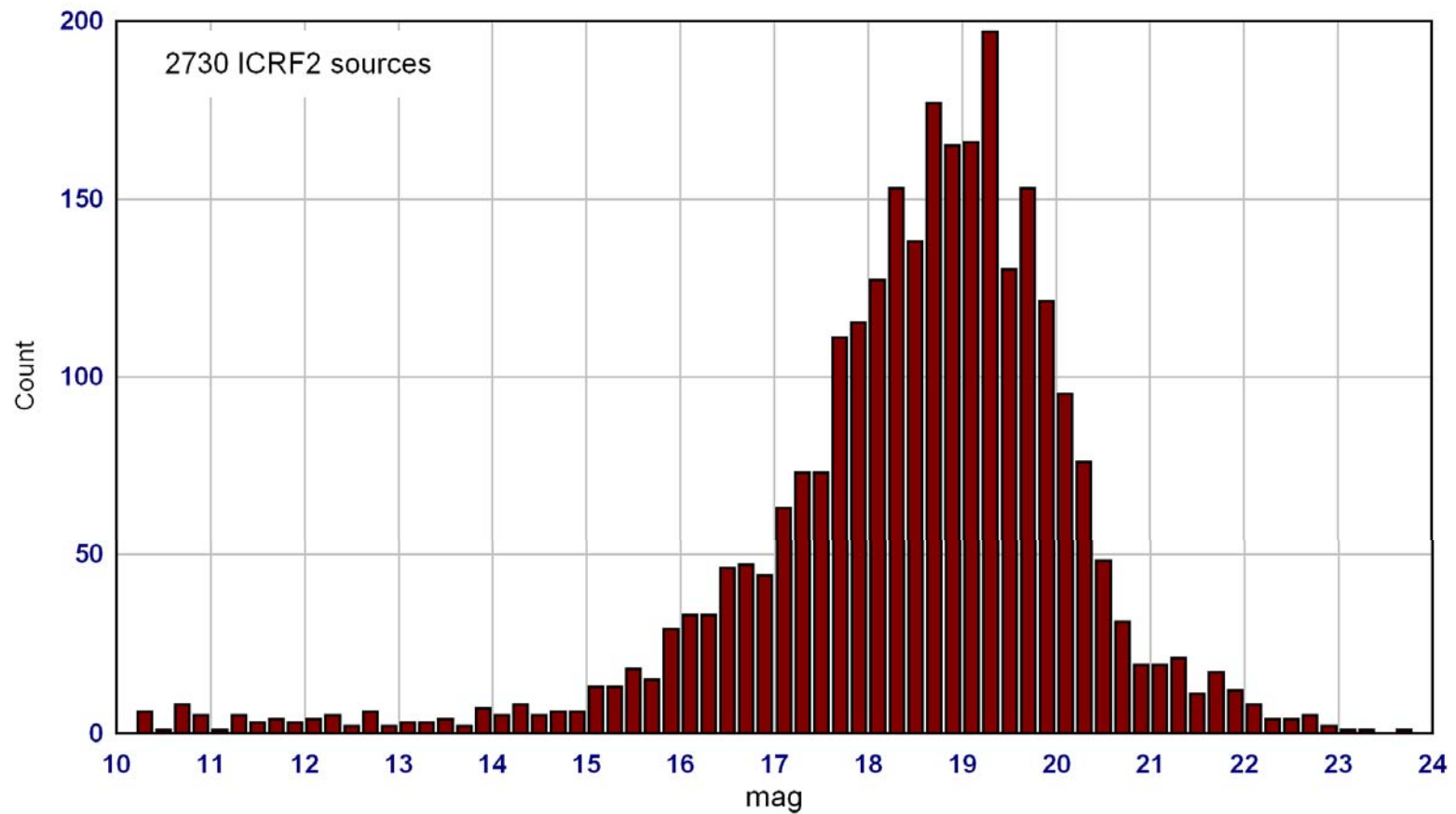
● VLBA Calib. (2197)

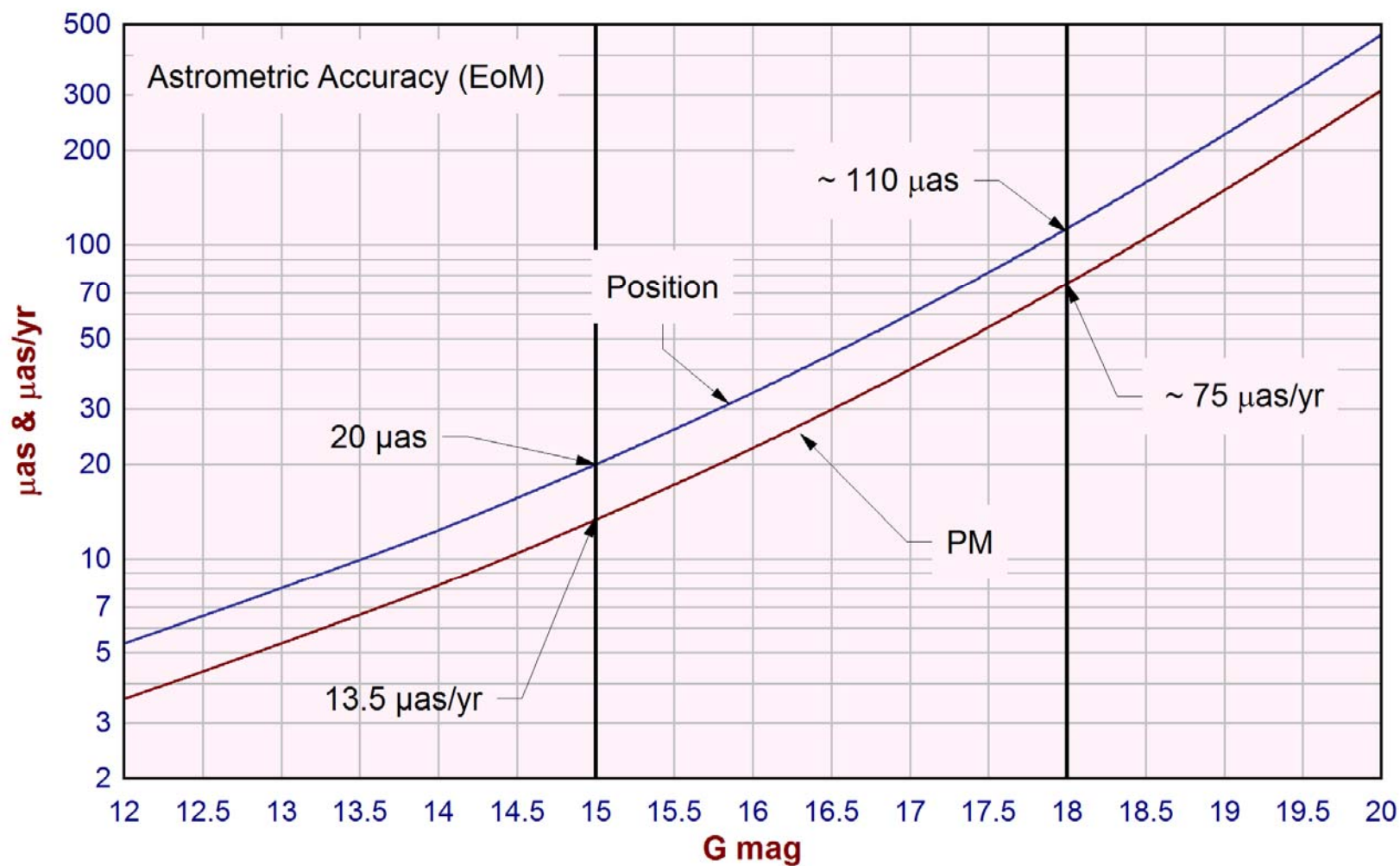
- Relatively faint sources in the visible
 - From ICRF1 B Magnitude given by Veron & Veron-Cetty
 - $V \sim 17-21$
- No visual magnitude available in the ICRF2 publication
- Cross-matched of the LQAC with ICRF2
 - For each entry there are between 0 and 7 photometric bands available

#	nom_source	ra	dec	flag_cross	u	b	v	g	r	i
	LQAC_000-032_001	0.084999781	-32.350342643	AB-----I-KLM	0.00	18.57	17.00	0.00	17.99	17.86
	LQAC_000+040_001	0.221173199	40.900498078	AB-D-----	0.00	0.00	0.00	0.00	0.00	0.00
	LQAC_017-060_001	17.314480025	-60.830127769	AB-----	0.00	0.00	0.00	0.00	0.00	0.00

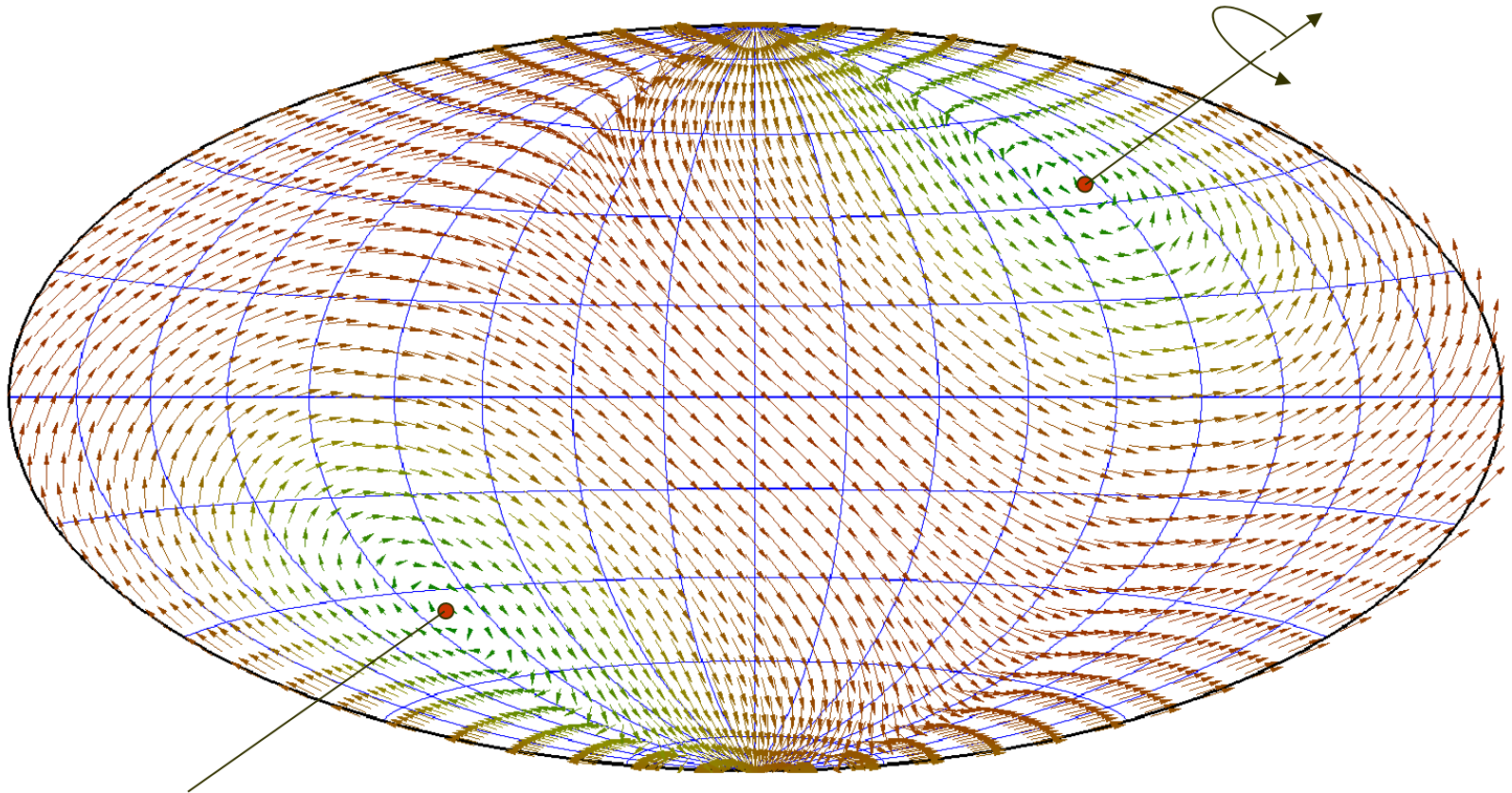
- When > 2 bands available, G magnitude can be estimated
 - otherwise use V or R as substitute
 - or systematic shift if only u, b are available
- Out of 3400 ICRF2 sources, I ended up with 2700+ with an estimated G magnitude

- Magnitude distribution

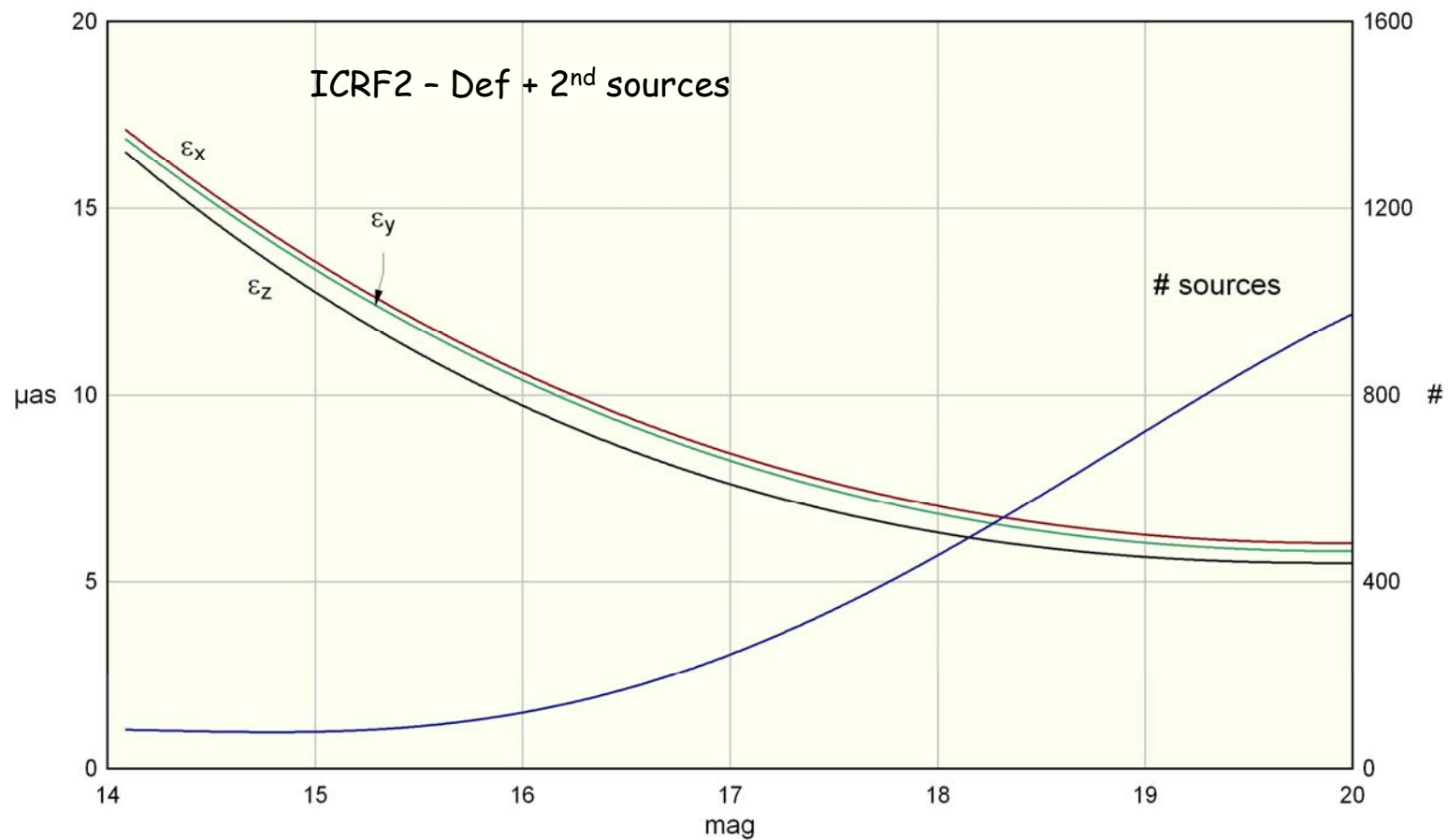




- Pattern solved with three parameters ($\varepsilon_x, \varepsilon_y, \varepsilon_z$)
- same for proper motion with ($\omega_x, \omega_y, \omega_z$)

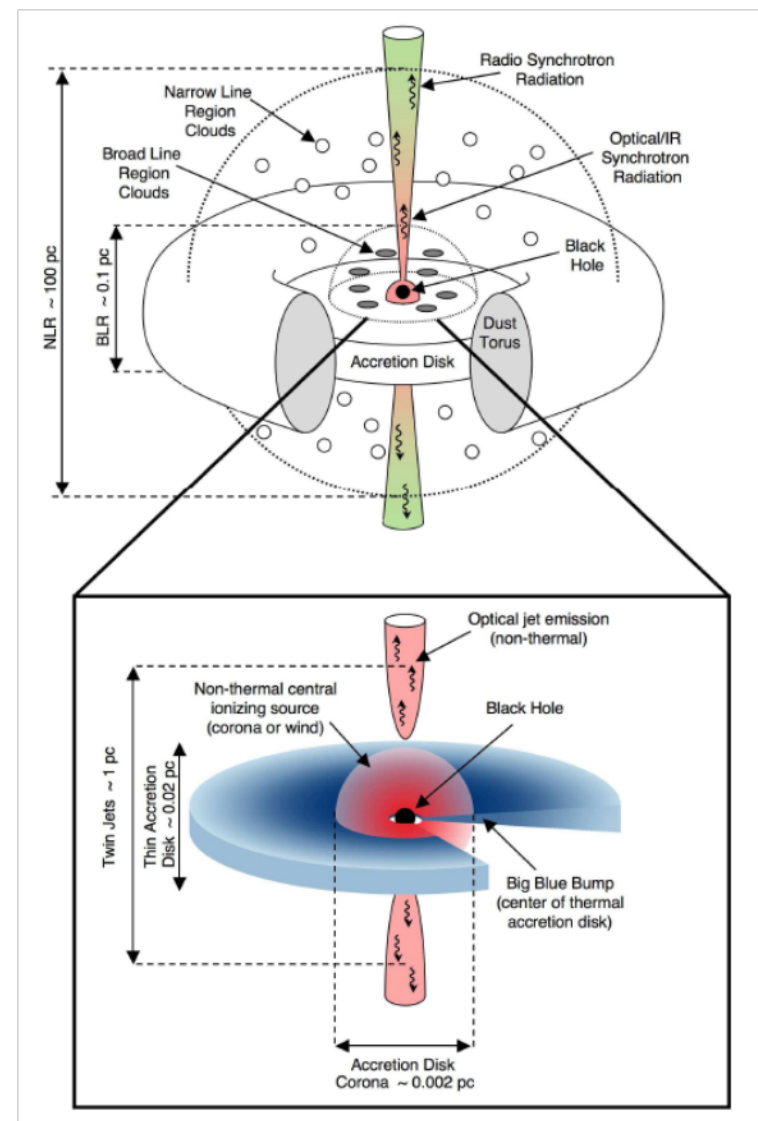


- based on the covariance matrix of the rotation model

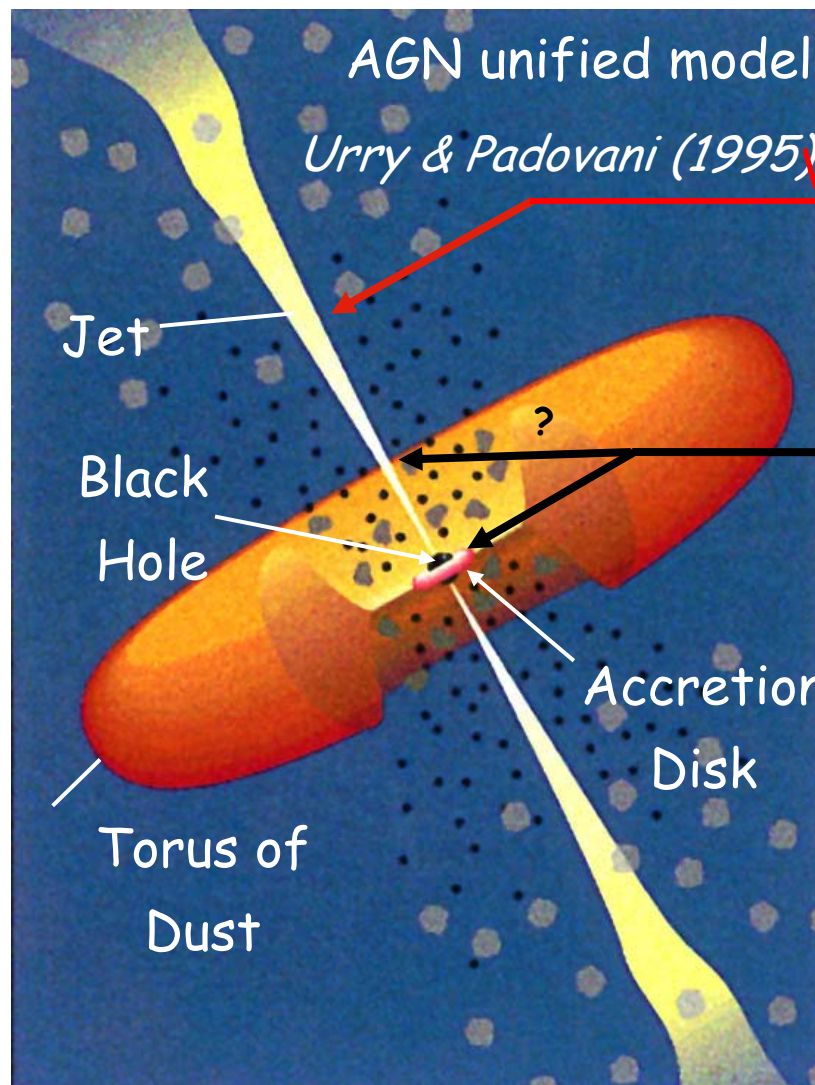


- The number of suitable sources for the alignment may be small
 - stable, no structure, bright enough in the visible
 - About 200 ICRF2 sources are suitable (Bourda et al., 2012)
- A program of selection and observation of additional sources is under way
 - PI : G. Bourda (Obs. of Bordeaux)
- Initially 450 sources pre-selected
 - Flux > 20 MJy, $\delta > -10^\circ$, $V < 18$
 - 400 found detectable in VLBI
- Imaging of 250 sources
 - 120 found without structure
- Astrometric observations of these sources in progress
- We may have > 120 new and high-quality source for the alignment
- ICRF3 will be also an important data set to explore

- No observational evidence
 - offset < 0.1 to 1 mas
- Emission model can provide insights
- For a distance of 0.1 pc between photocenter and radiocenter \rightarrow angular distance $> 10 \mu\text{as}$ ($z > 1$, RG involved)
 - becomes relevant for the alignment on radio ICRF
 - the offsets vectors should be randomly oriented \rightarrow act as an additional noise



Frequencies
in VLBI:
S ~ 2 GHz
X ~ 8 GHz
K ~ 24 GHz
Ka ~ 32 GHz
Q ~ 43 GHz



VLBI observation

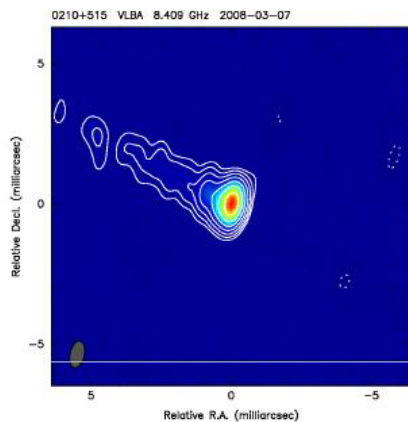
~ 100 μ as
Kovalev et al. 2008

Gaia
observation

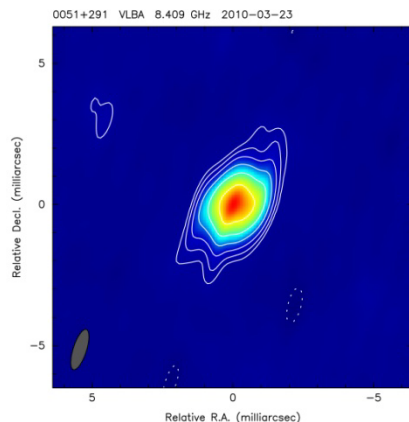
credit : G. Bourda, Obs. Bordeaux

VLBI maps of 'bad sources'

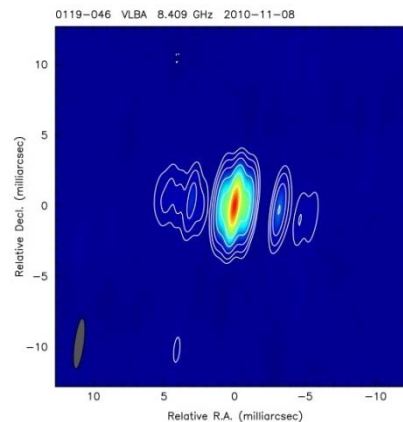
GC030



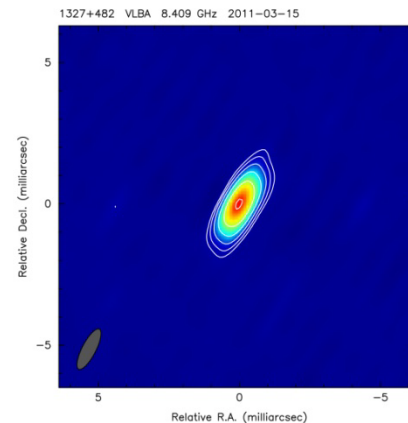
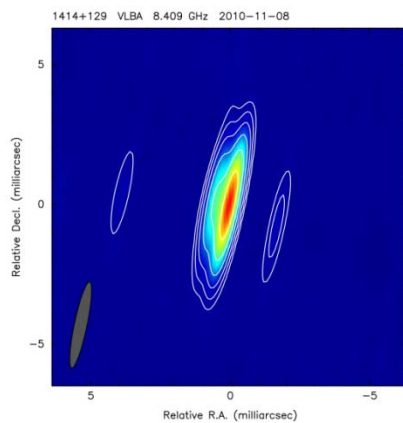
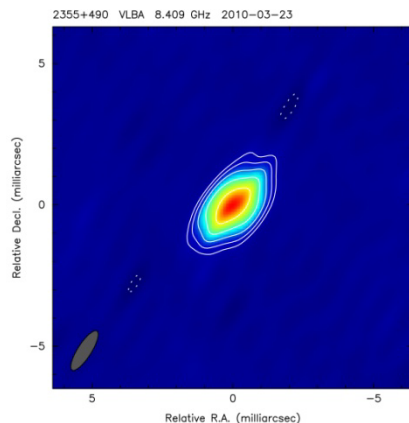
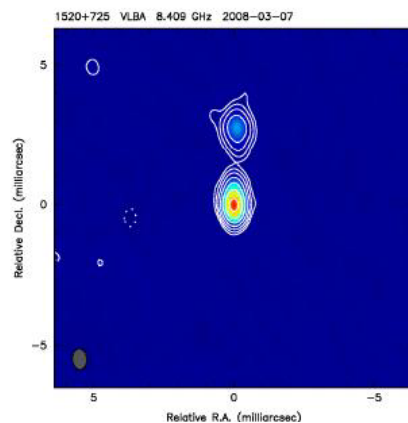
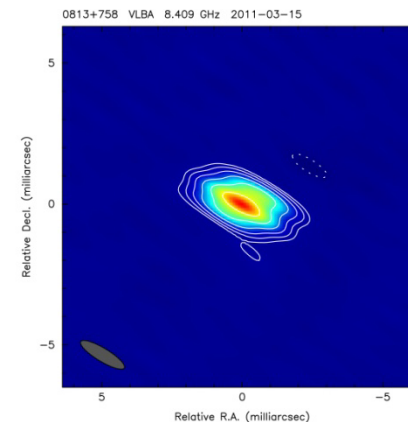
GC034A



GC034BC



GC034EF



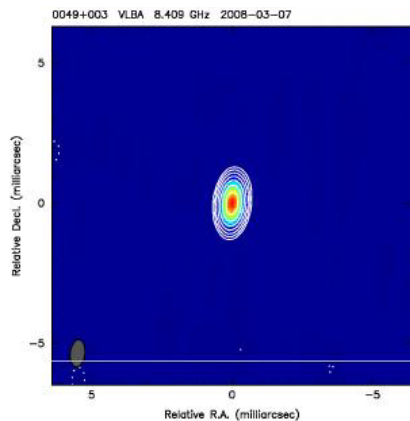
~10 mas

X-band -1st contour level @ 1 - 4%

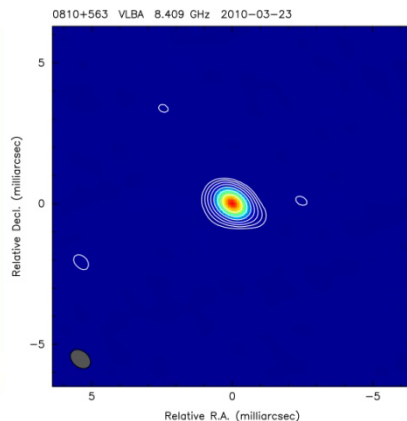
credit : G. Bourda, Obs. Bordeaux

Santiago, 3/10/2014 - F. Mignard

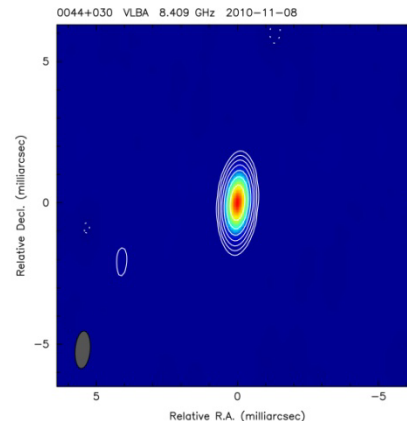
GC030



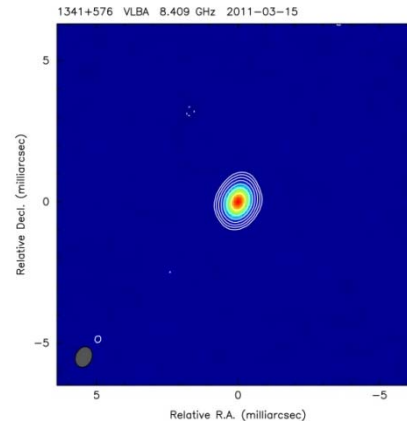
GC034A



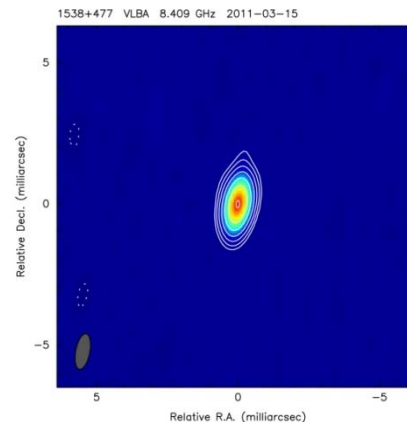
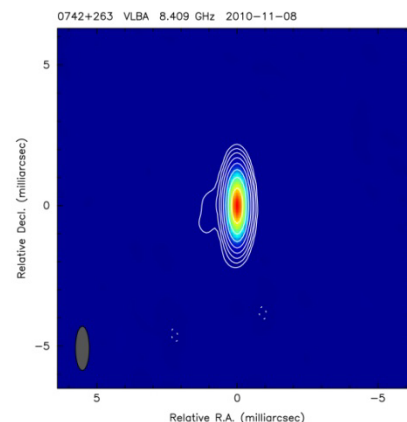
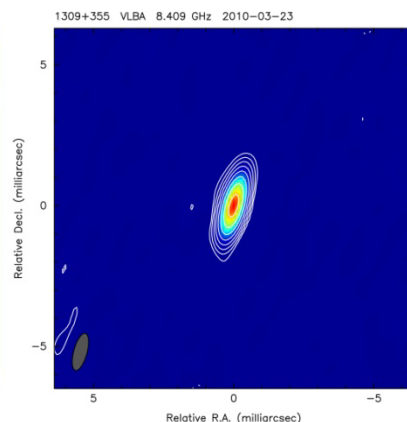
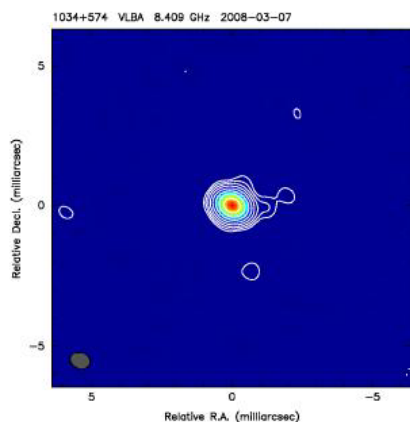
GC034BC



GC034EF



~10 mas



X-band -1st contour level @ 1 - 4%

credit : G. Bourda, Obs. Bordeaux

- No split between rotation and cosmic acceleration
- Orientation must be determined at the same time
 - cosmic proper motions to be included in the ICRF positions
 - they are not known, but can be computed
- Needs to carry out the alignment:
 - small subset of sources with positions in ICRS
 - ICRF sources and other observed in VLBI → alignment
 - larger subset of EGSs with statistically zero proper motions
 - not necessarily observed with VLBI → rotation

- There are 9 reference frame parameters
 - 3 for the orientation $\varepsilon_x, \varepsilon_y, \varepsilon_z$
 - 3 for the rotation $\omega_x, \omega_y, \omega_z$
 - 3 for the cosmic acceleration a_x, a_y, a_z
- They are all determined within the astrometric global solution
 - attitude, position, proper motion will be referred to this frame
 - hopefully the ephemeris is given in a frame nearly identical
 - however the coupling is weak, and an error of 10 mas is acceptable
- This should be carried out in post-astrometric processing